

# Quarkonium Production in Hot Medium

Ágnes Mócsy

Pratt Institute, Brooklyn, New York

**Pratt**

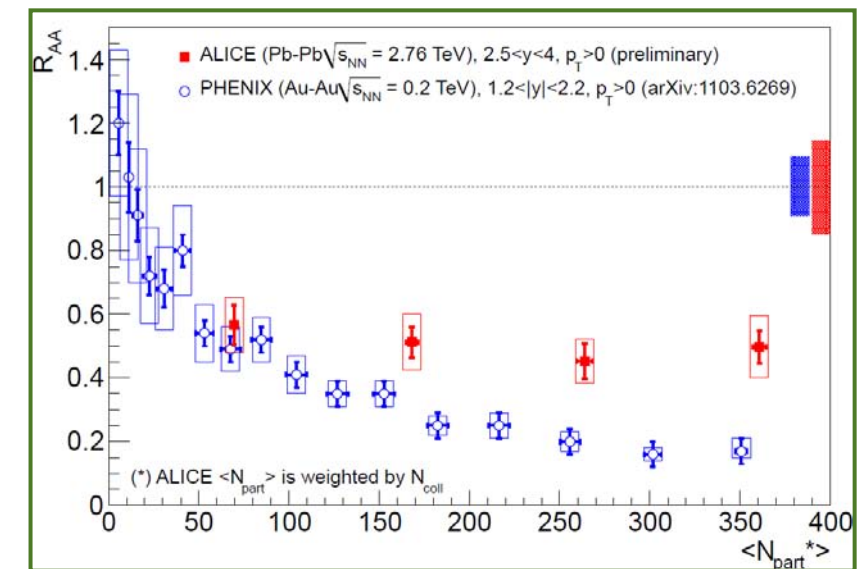
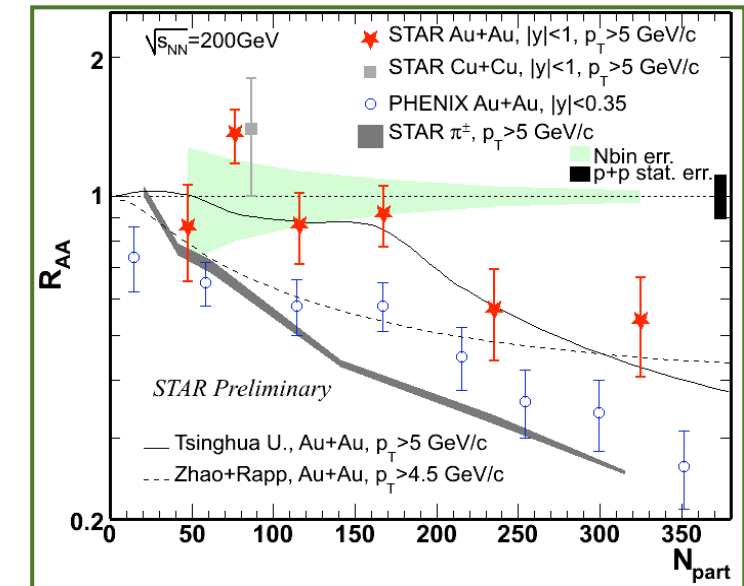
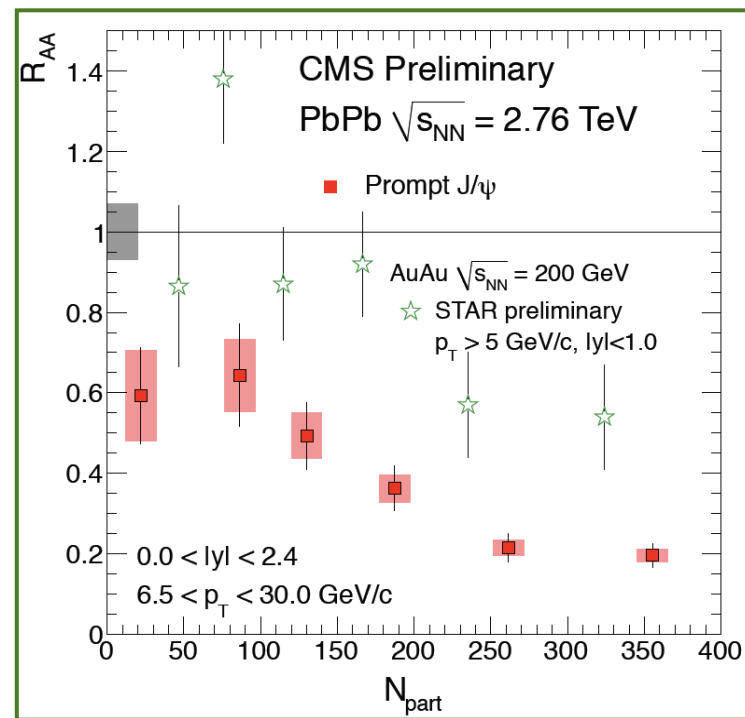
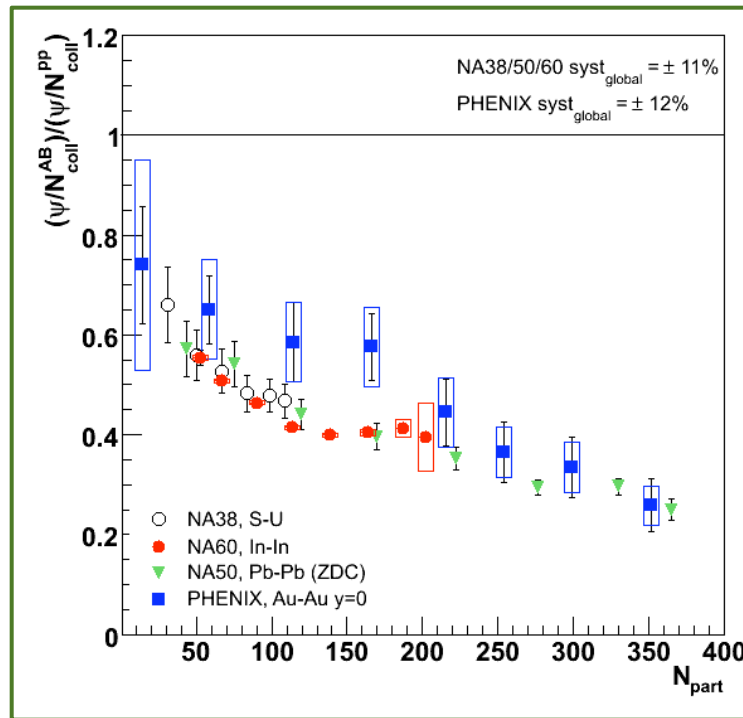


# Herzliche Glückwünsche, Helmut!



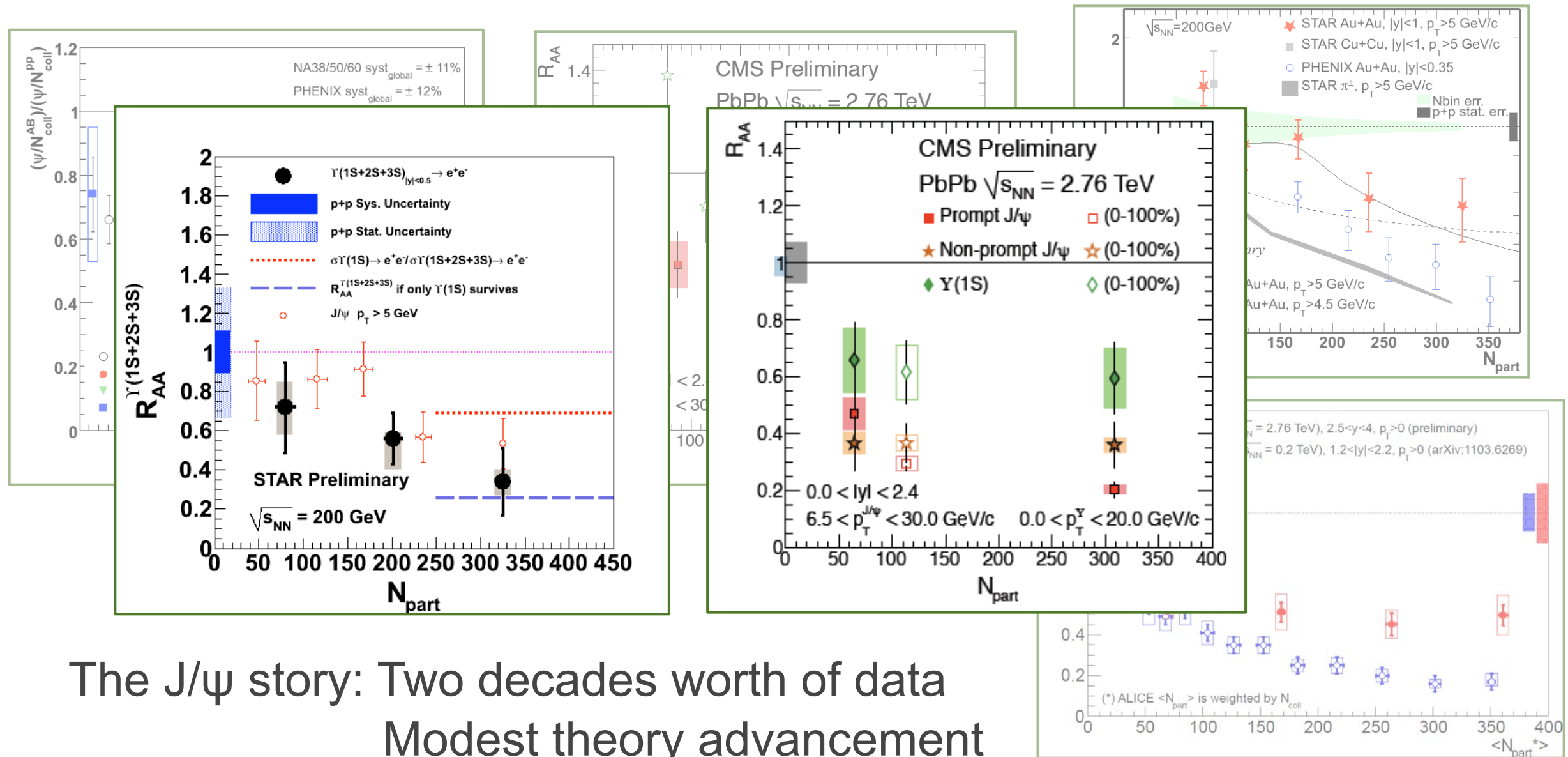


# Starting Point : $R_{AA}$



The J/ $\psi$  story: Two decades worth of data  
Modest theory advancement  
Lots of ad-hoc phenomenological modeling  
It is difficult to unambiguously interpret - we are still not there

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The  $\Upsilon$  story just started !

# Main Outline

- What we know theoretically about quarkonium in deconfined medium
- Bridging between theory and experimental data
- Some cross-checks: in attempt to isolate pure hot medium effects

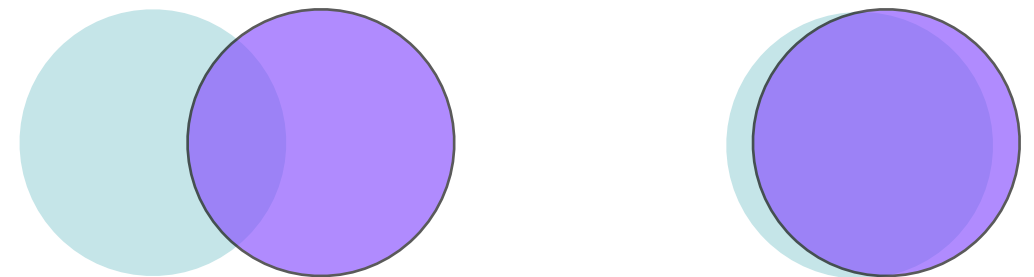
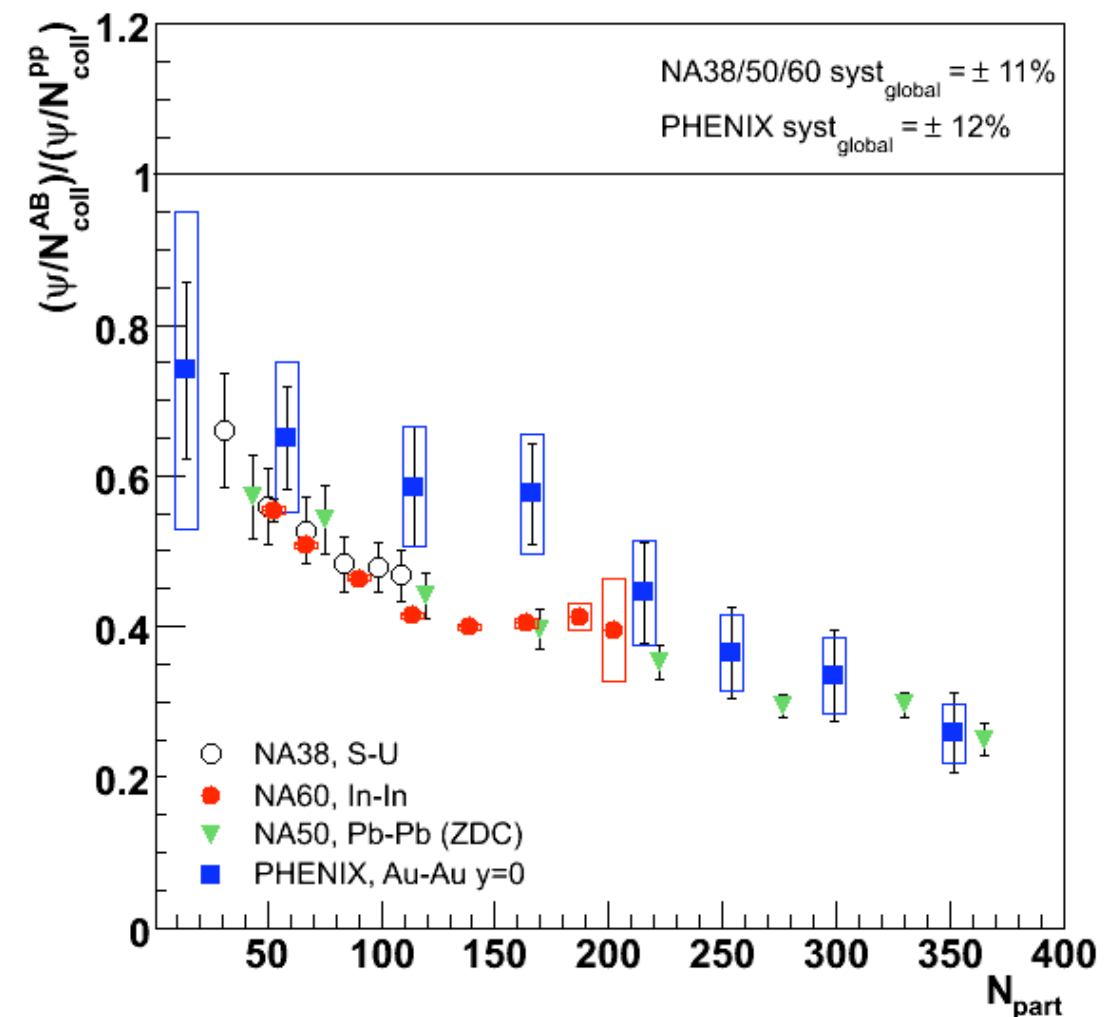
# The $R_{AA}$

- $J/\psi$  nuclear modification factor:  
yield in AA collisions relative to yield in pp  
(where no QGP formation expected) scaled with  
number of binary NN collisions

$$R_{AA}^{J/\psi} = \frac{dN_{J/\psi}^{AuAu}/dy}{N_{coll} \cdot dN_{J/\psi}^{pp}/dy}$$

- If AA is superposition of pp then  $R_{AA}=1$
- Deviation from 1 indicates medium effects
- If no  $J/\psi$  measured then  $R_{AA}=0$

- A  $J/\psi$ -suppression pattern observed at  
SPS and RHIC and LHC



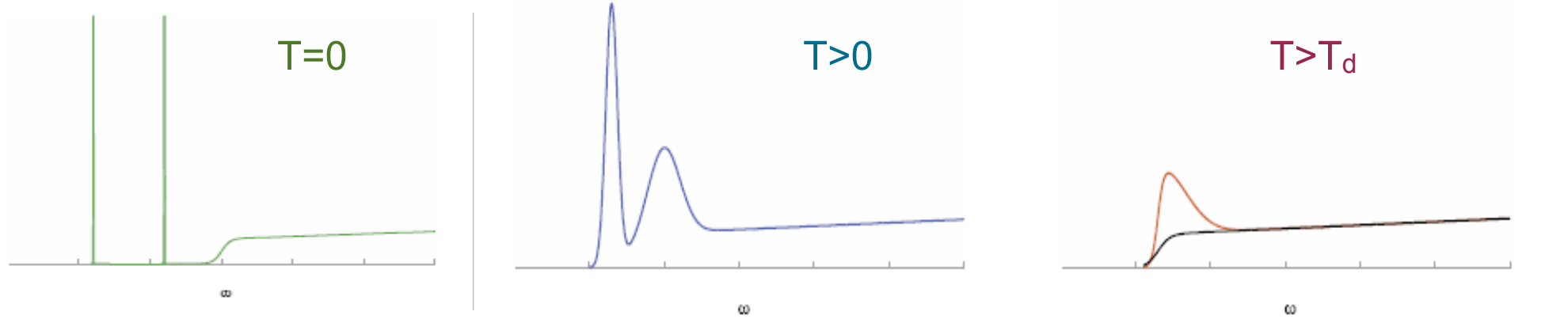
# QCD Expectations

- QCD predicts: quarkonium states disappear in the deconfined medium due to temperature effects present: *screening, Landau damping, ...*

- In-medium properties of quarkonium encoded in spectral functions

$$\sigma(\omega, p, T) = \frac{1}{2\pi} \text{Im} \int_{-\infty}^{\infty} dt e^{i\omega t} \int d^3x e^{ipx} \langle [J(x, t), J(x, 0)] \rangle_T$$

Dissolution (“melting”) seen as progressive broadening and disappearance of bound-state peaks

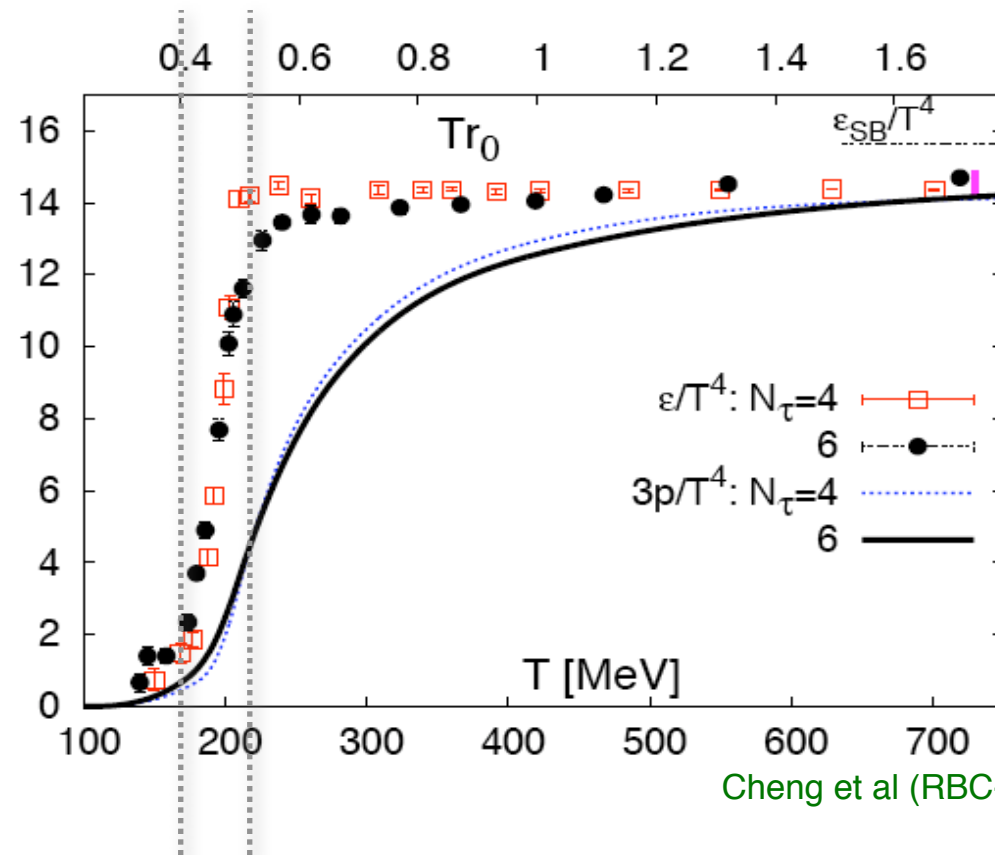


- Theory predicts the J/ψ disappears in the plasma → J/ψ suppression proposed signal of deconfined QGP

Matsui, Satz PLB 1986

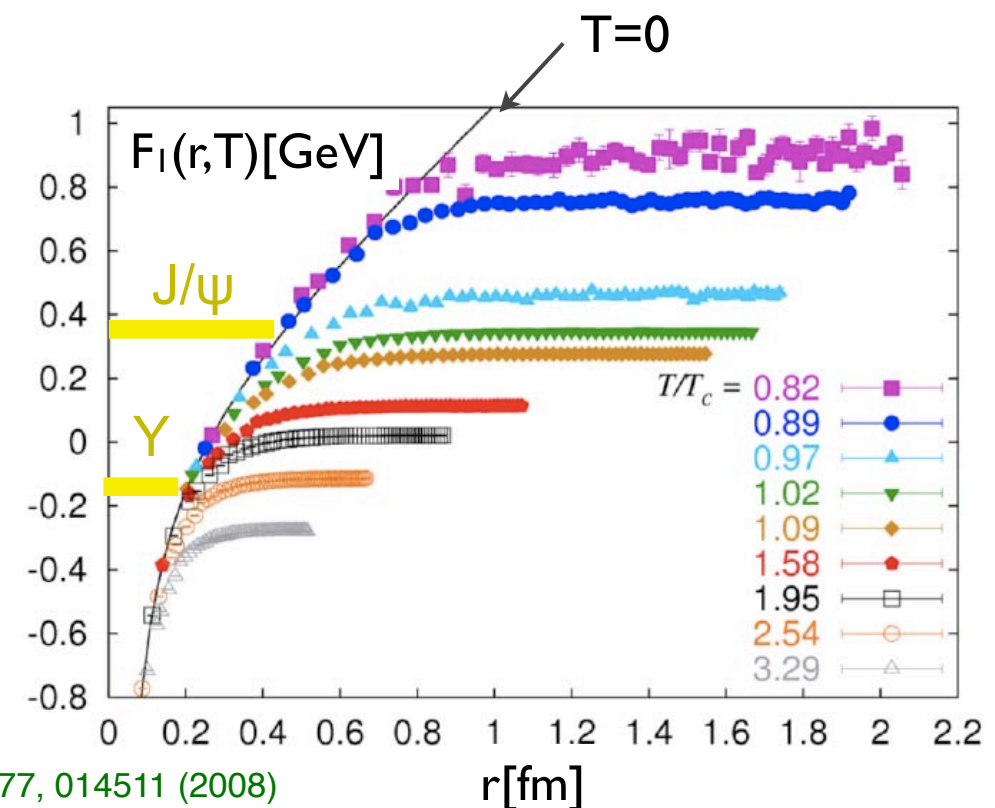
# What Do We Know Now

## Deconfinement



- Rapid rise of the energy density: liberation of new degrees of freedom
- Deconfinement seen on lattice

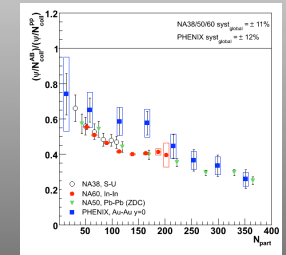
## Color screening



- Strong screening of static Q-Qbar free energy - sets in at shorter distances with increasing  $T$
- $r_{scr} < r_{J/\psi}$  “melting” of the  $J/\psi$



# What's the Physics Behind



But is the  $J/\psi$   $R_{AA}$  a signal for deconfinement and screening ?!

To answer this question we need to know:

- How the properties of  $J/\psi$  change in a deconfined medium  
Determine the spectral function
- Relate an equilibrium spectral function to  $R_{AA}$   
Through real-time dynamics
- Identify what physics might contribute to  $R_{AA}$   
for example: Suppression is seen in pA, dA data as well  
(where no QGP formation expected)  
Cold nuclear matter effects - could be relevant to AA

# Theory Progress

- Lattice QCD, potential models, effective field theories (EFT)

- **Spectral functions are calculated**

Cabrera, Rapp, Mócsy, Petreczky, Alberico, Beraudo, ...

- Potential model assumes most medium effects on quarkonium properties can be described by a T-dependent potential

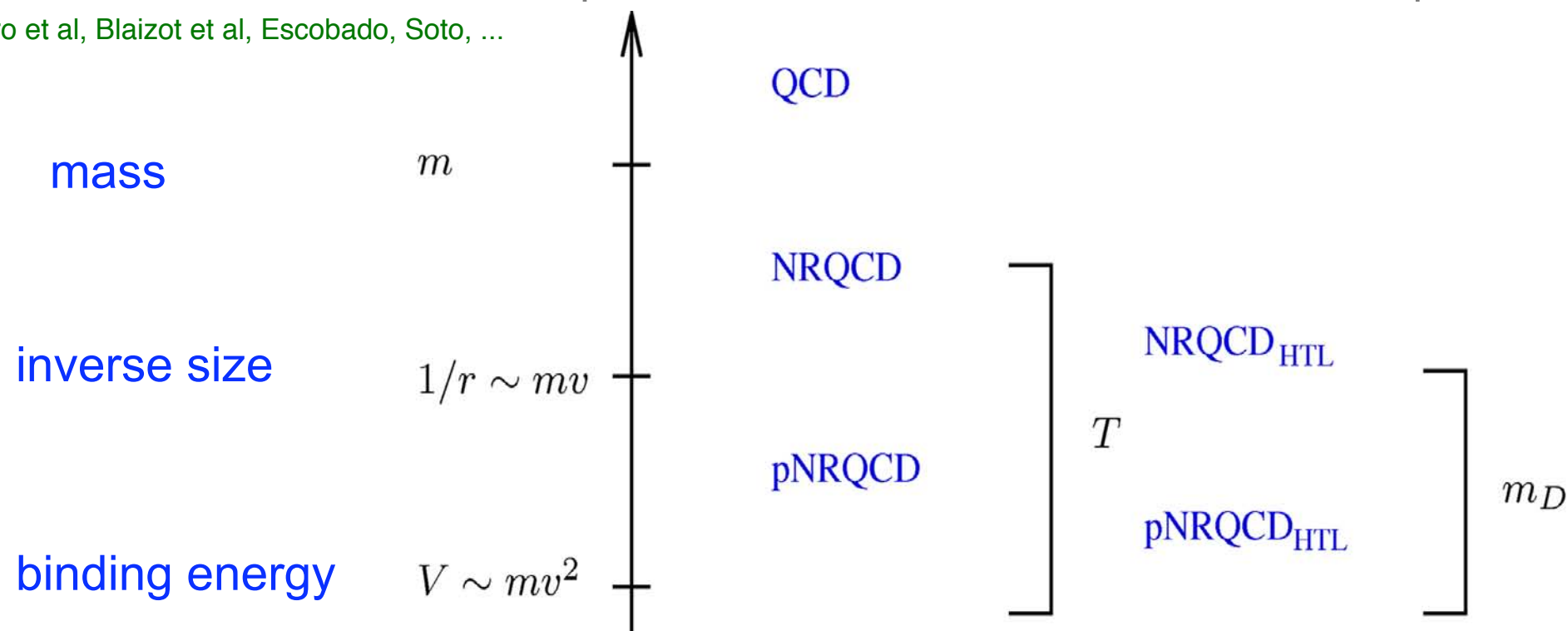
where  $E_{\text{bin}}$  is the smallest scale

- Recently: Potential model (can be placed on more solid grounds) appears as the tree-level approximation of the EFT and can be systematically improved

The heavy quark mass provides a hierarchy of different energy scales

Scale separation allows to construct sequence of effective field theories: NRQCD, pNRQCD

Laine et al, Vairo et al, Blaizot et al, Escobado, Soto, ...



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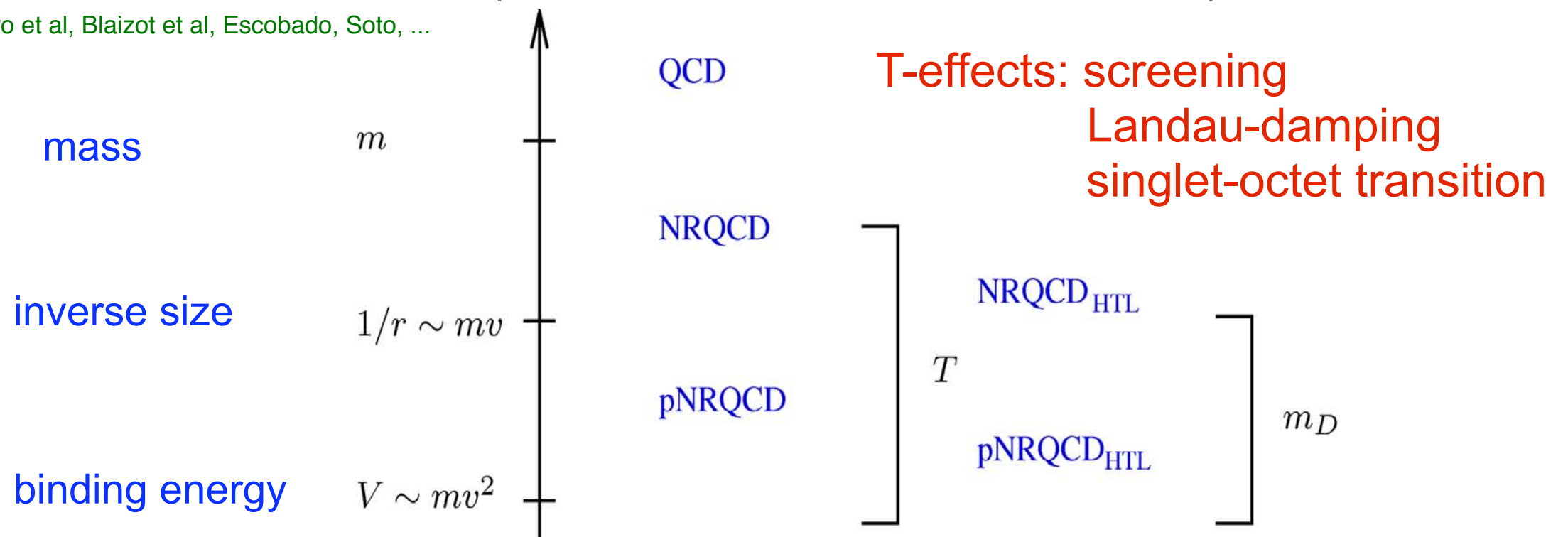
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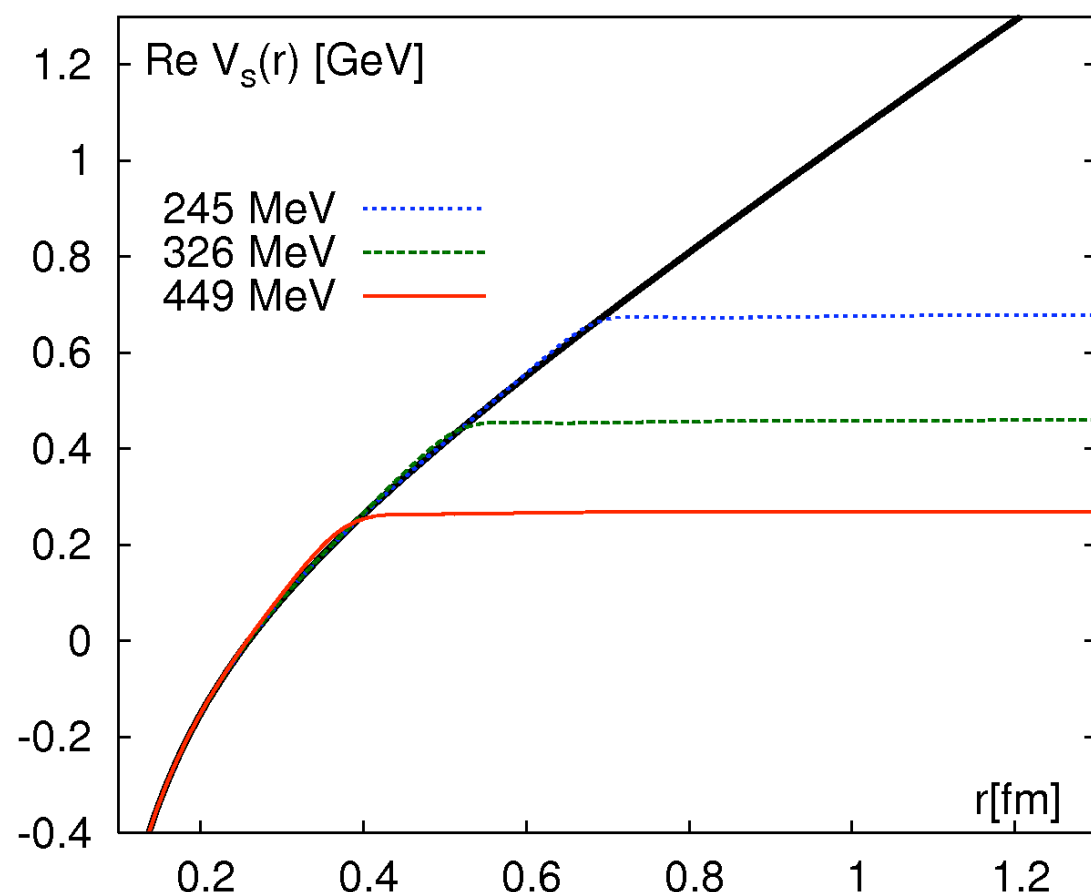


# The Complex Potential

Thermal contributions come from **Re** ( $\rightarrow$  binding energy) and **Im** ( $\rightarrow$  width) part

Constrain  $\text{Re}V_s(r)$  by lattice QCD data on the singlet free energy

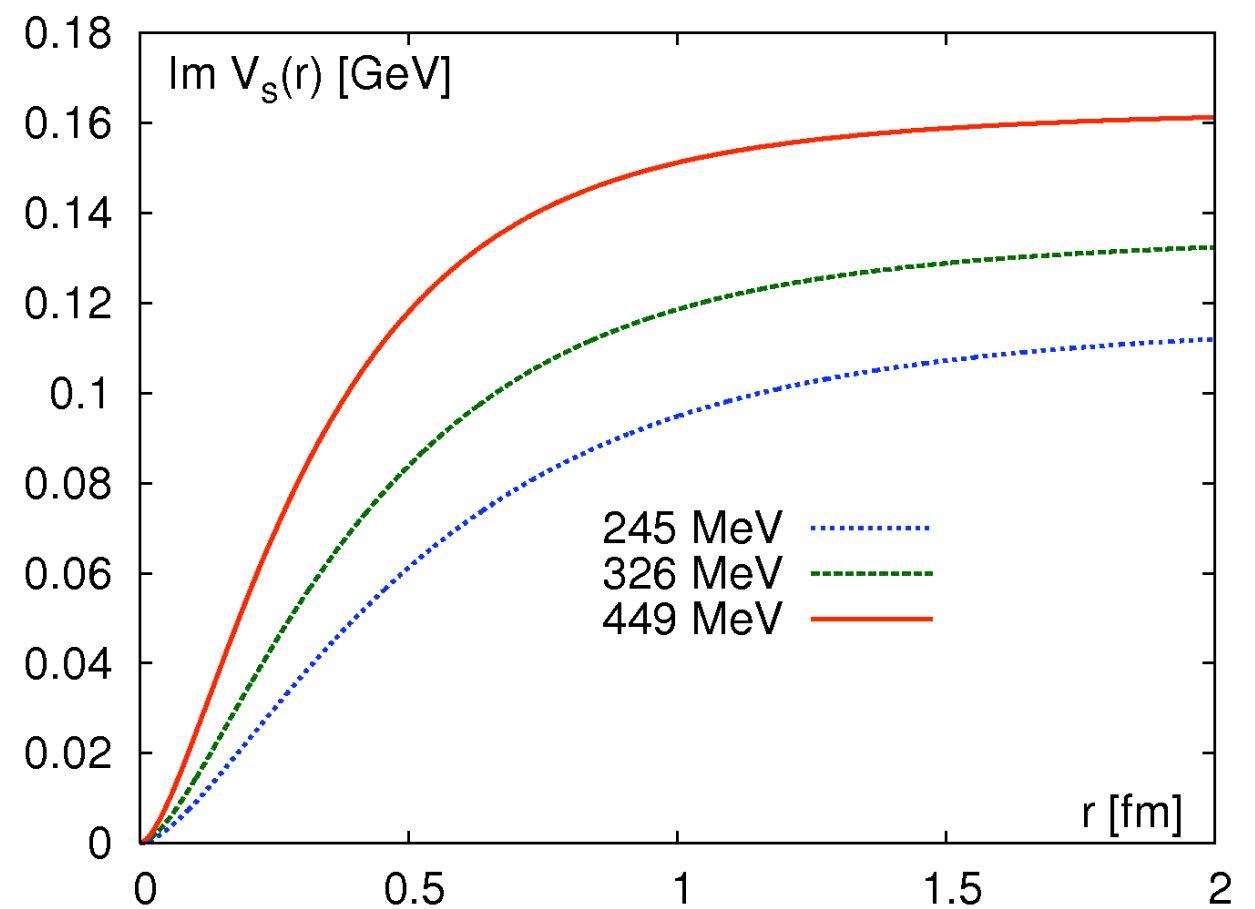
*Maximal value*



Mócsy, Petreczky, PRL 99 (07) 211602

Take  $\text{Im}V_s(r)$  from pQCD calculations

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Burnier, Laine, Vepsalainen JHEP 0801 (08) 043  
Beraudo, arXiv:0812.1130

Study the effect of **color screening** and of **dissipation** on the quarkonium spectral functions.



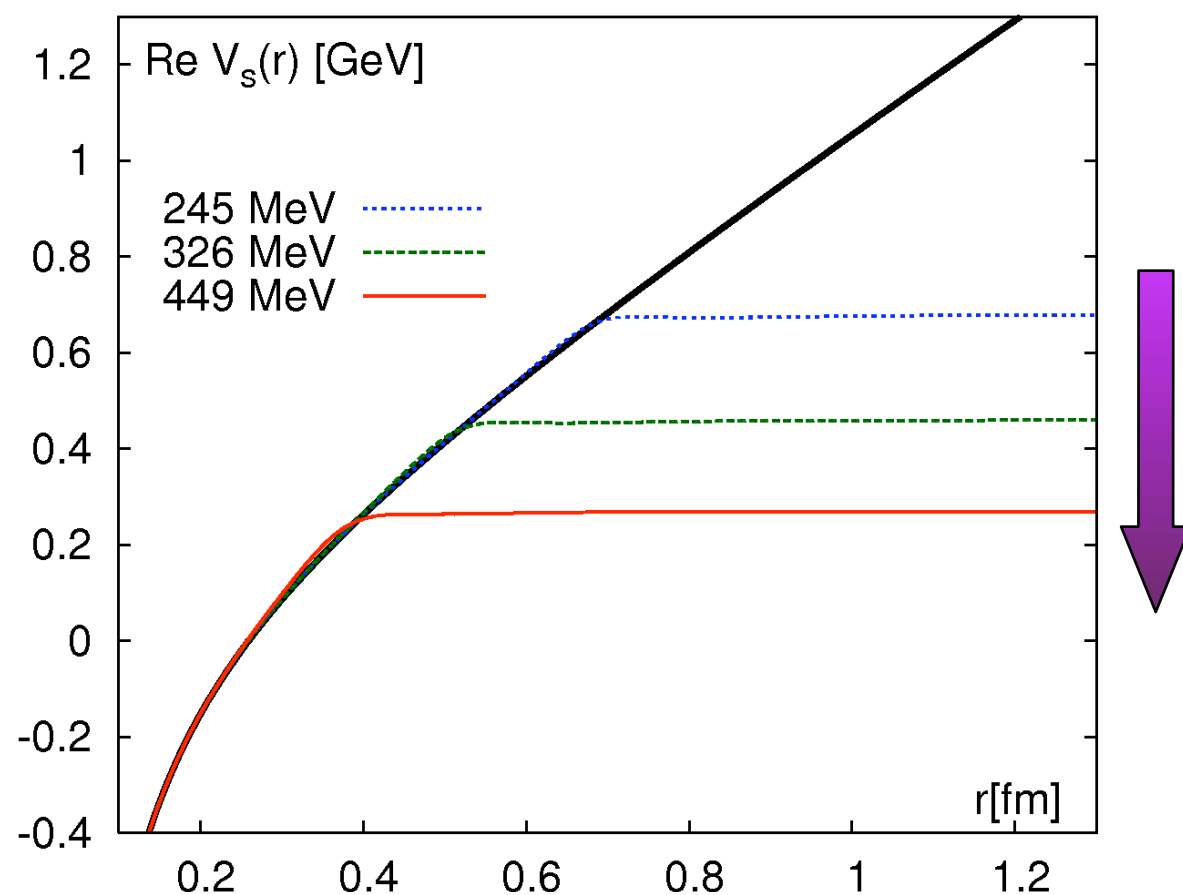
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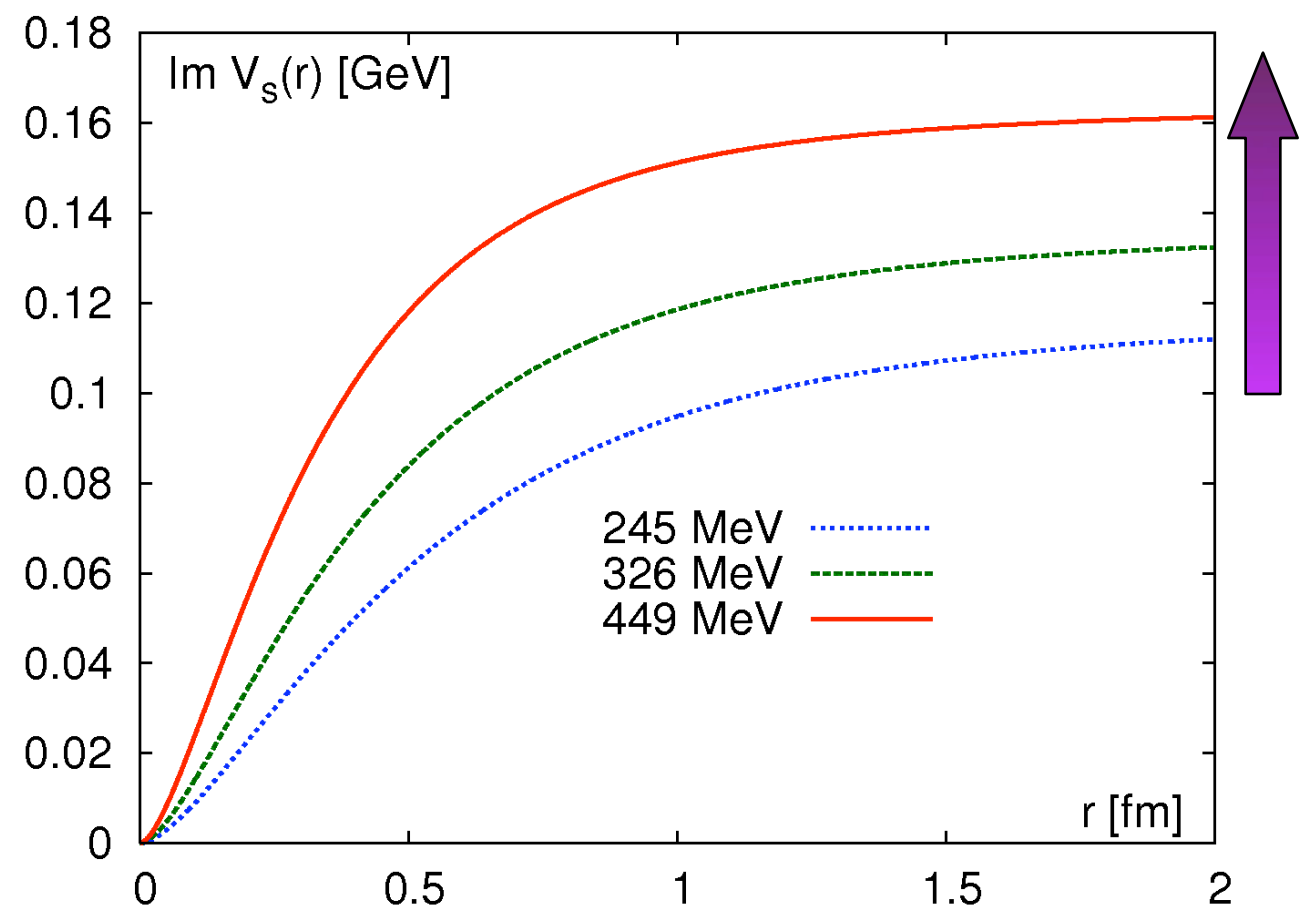
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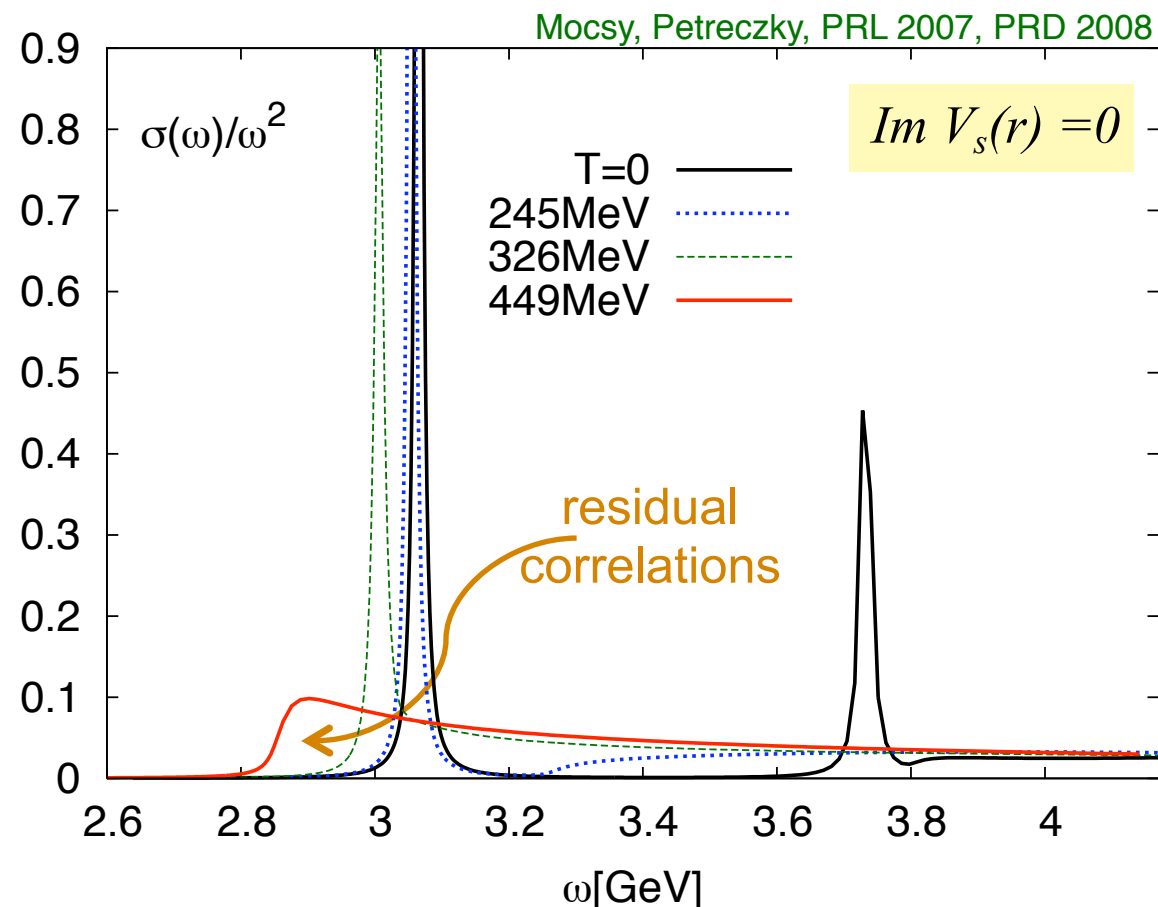
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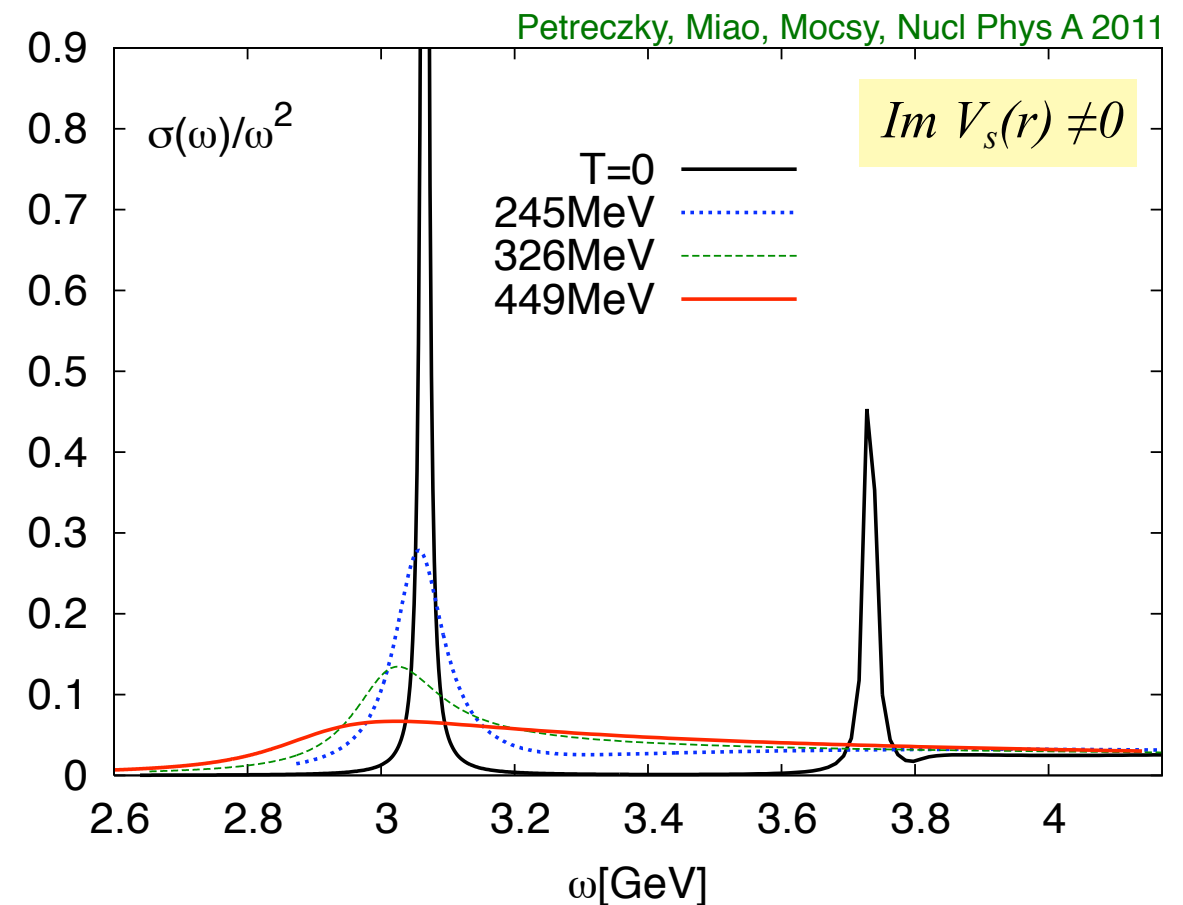
Calculated in full QCD

Take the **upper limit** for the real part of the potential allowed by lattice calculations



- 1S peak there with **binding energy reduced**
- residual c-cbar correlations persist  
threshold enhancement
- agrees with Cabrera, Rapp, PRD 2007

Take the perturbative imaginary part of the potential and the code from  
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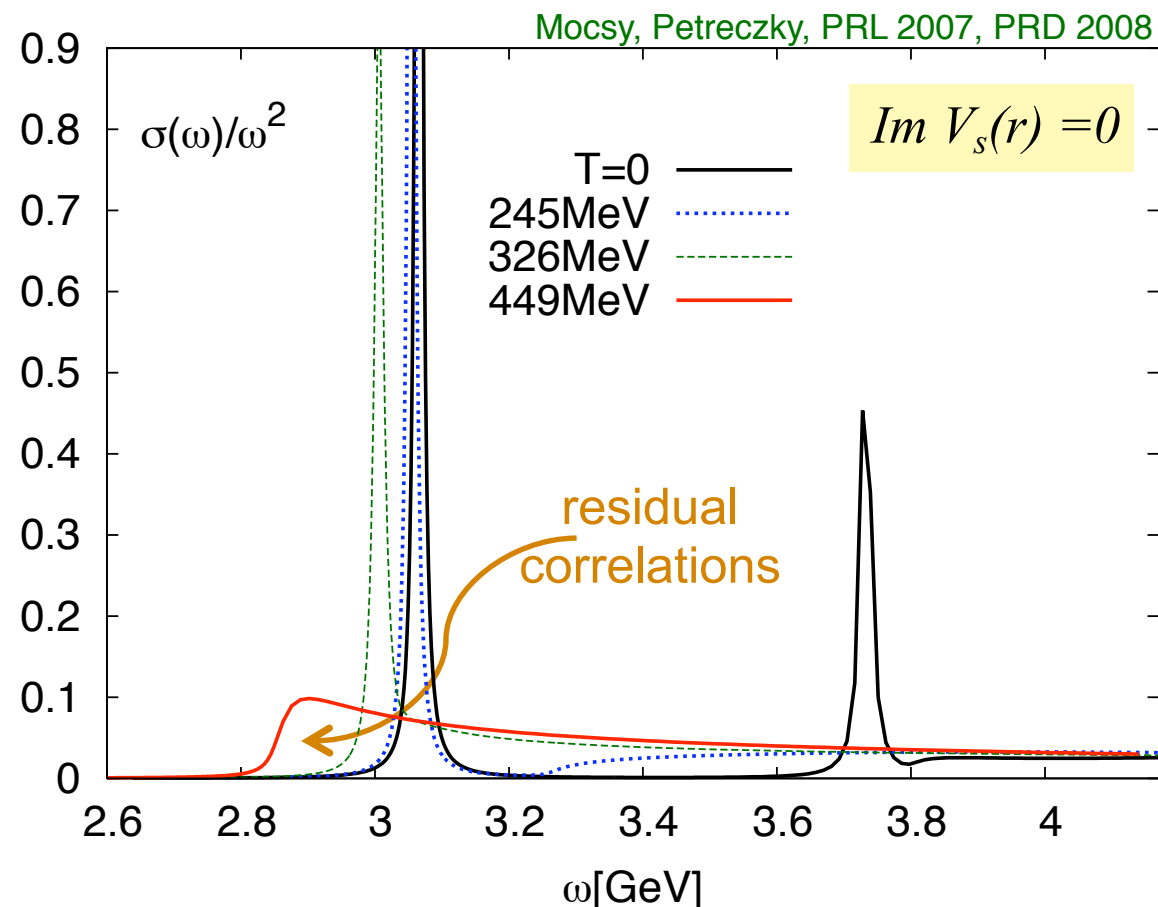


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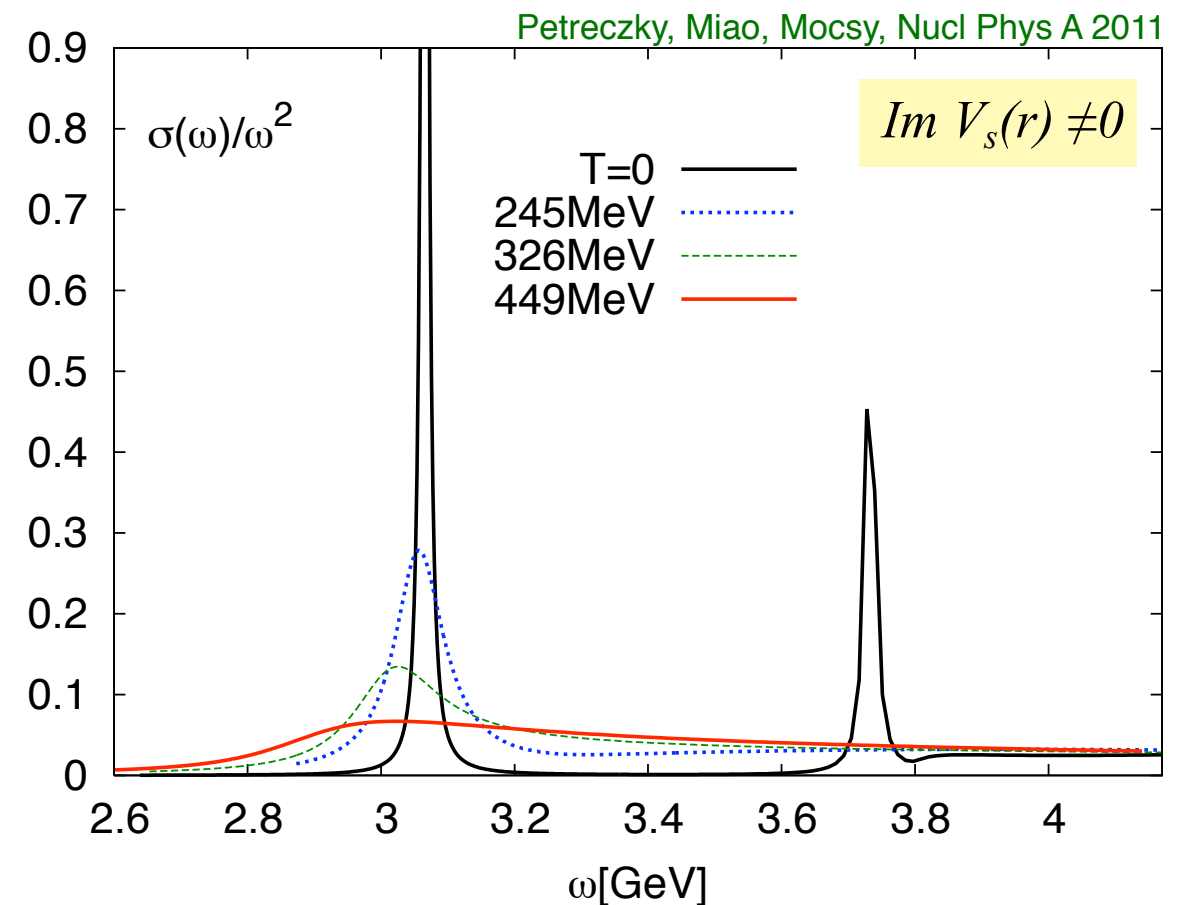
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**No charmonium state could survive for  $T > 240$  MeV**

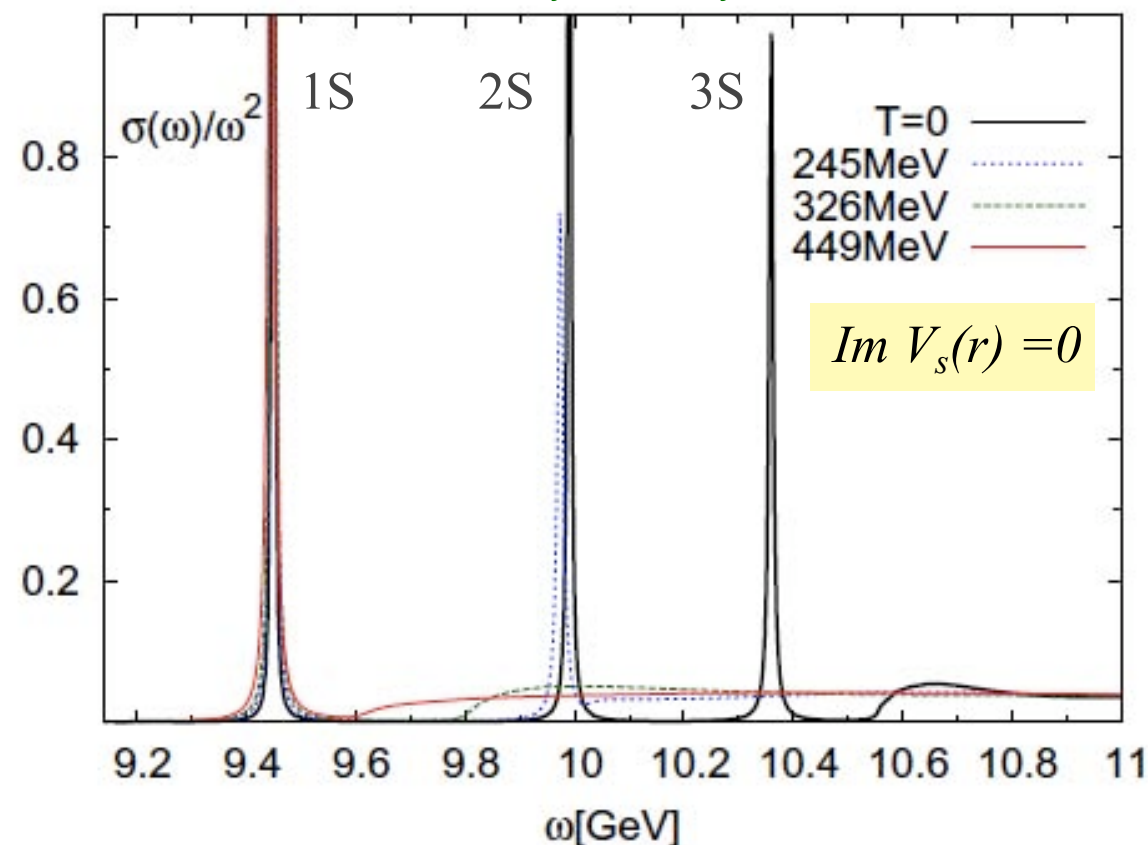
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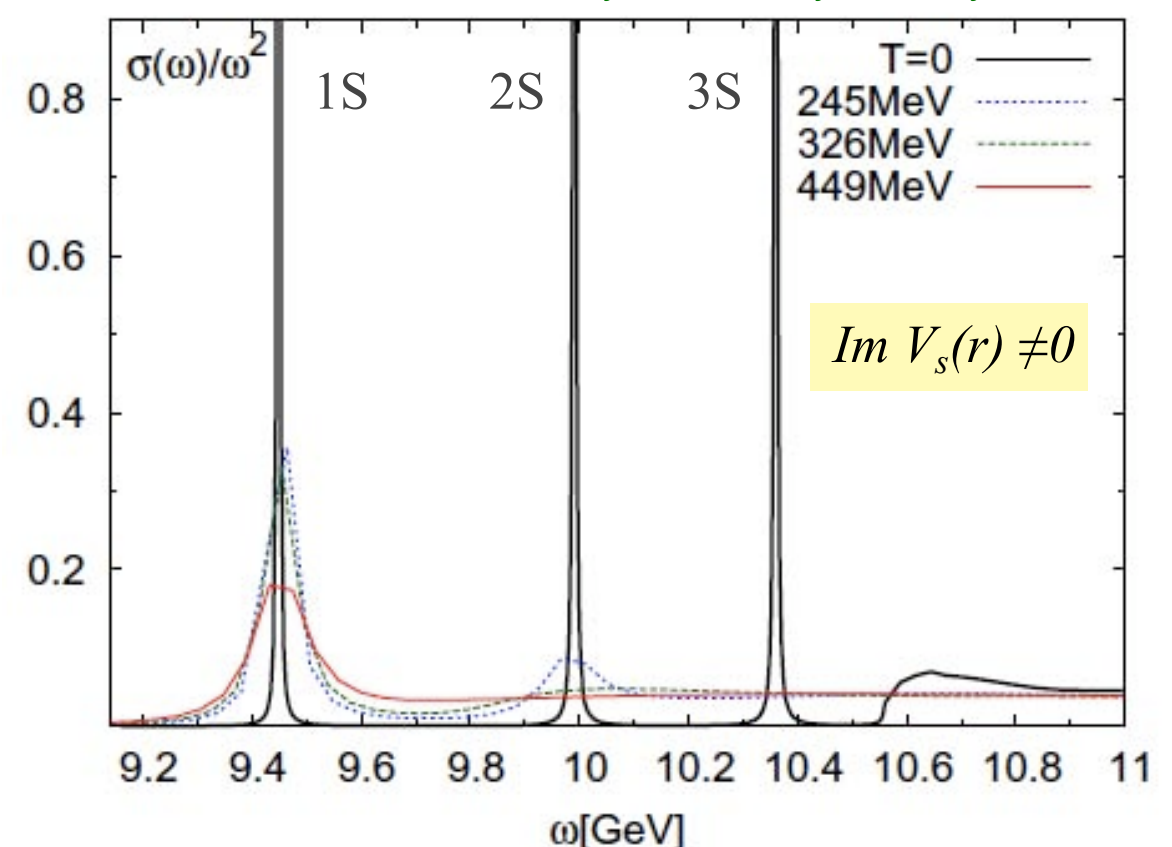
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Mocsy, Petreczky, PRL 2007, PRD 2008



- 1S peak and remnant of 2S state there
- binding energies reduced
- Threshold enhancement

Petreczky, Miao, Mocsy, Nucl Phys A 2011



- dramatic changes with Im part  
peaks significantly broaden



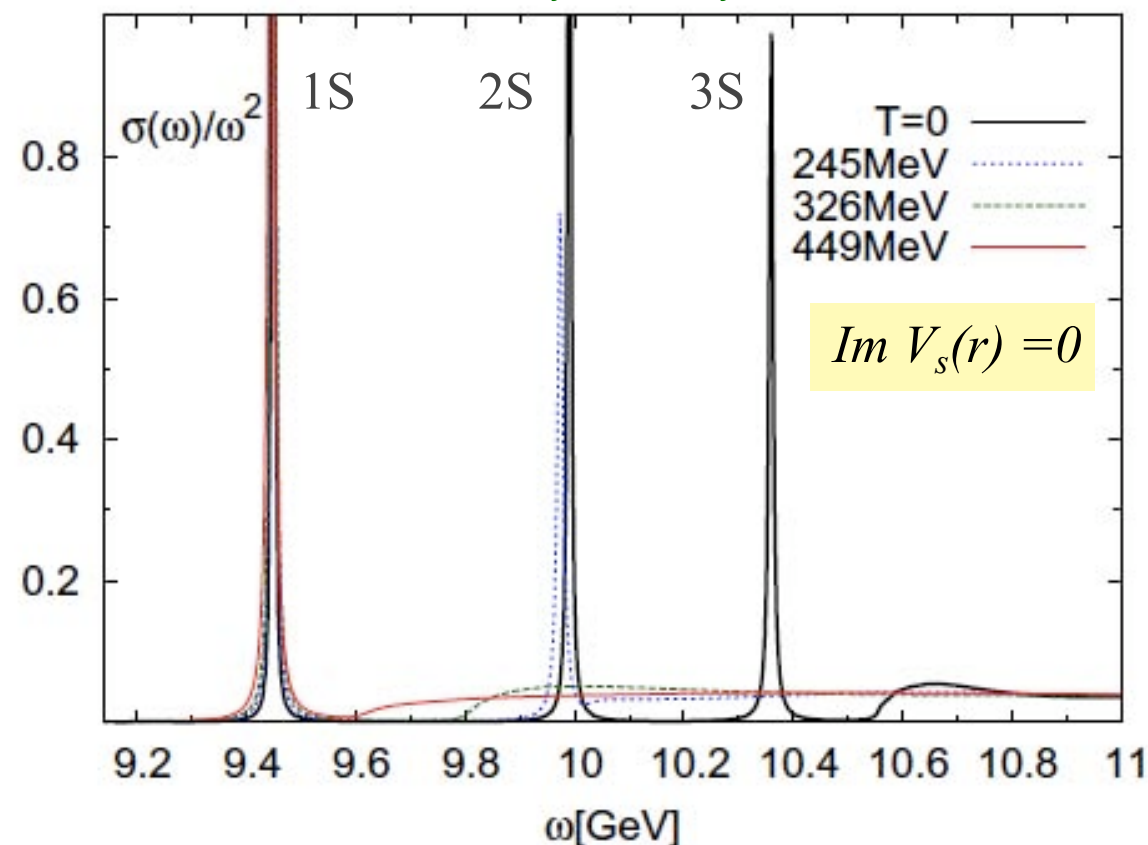
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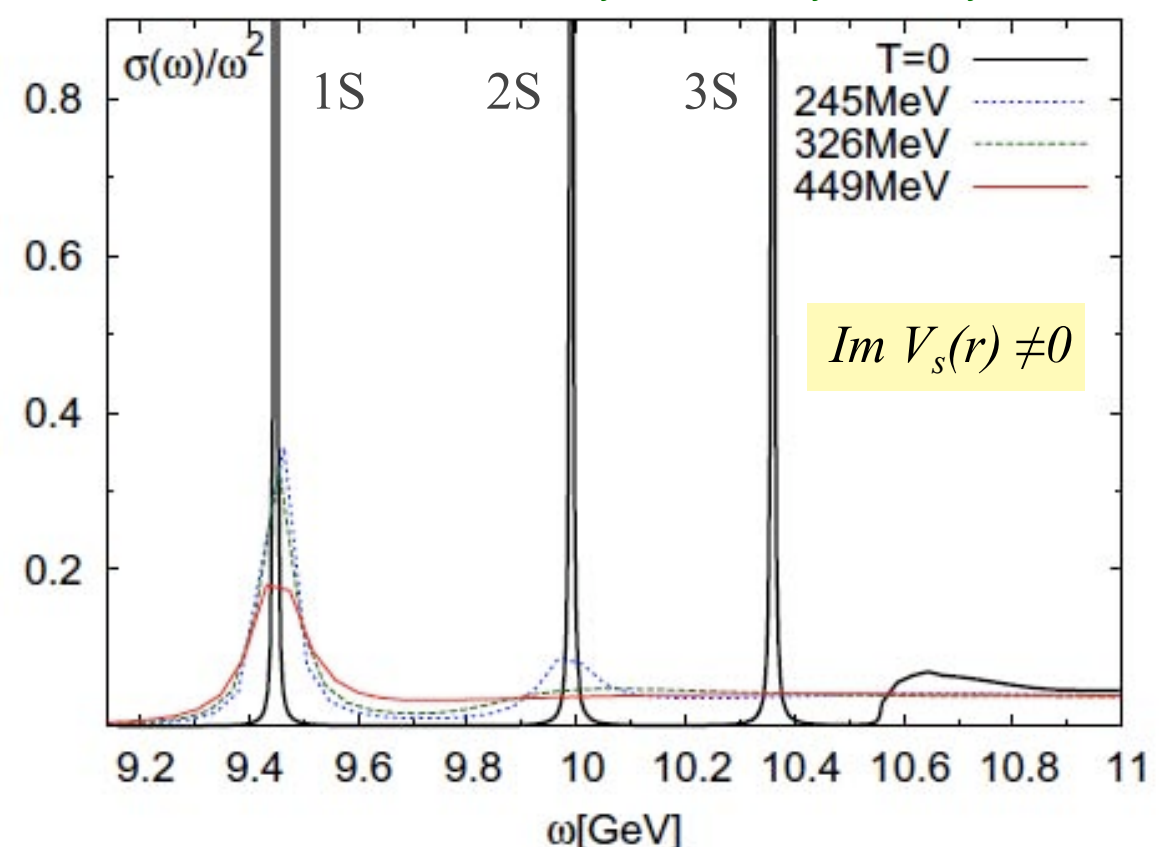
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**No bottomonium state could survive for  $T > 450$  MeV**

# Thermometer

- Quantitative estimates of peak disappearance

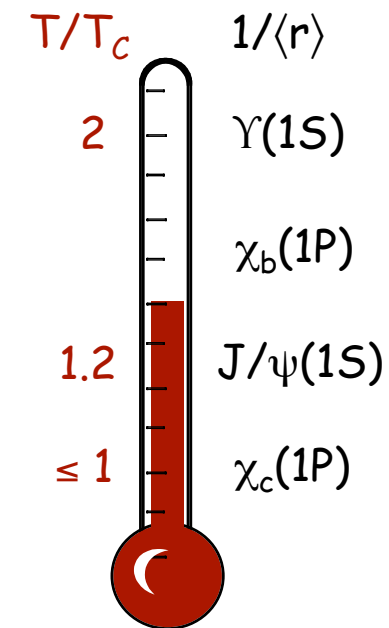
Thermometer of upper limits  $T_{\text{diss}}$

Note:  $E_{\text{bin}}$  would be smaller with other potentials

Charmonium sensitive to  $\text{Re}V$

Bottomonium sensitive to  $\text{Im} V$

Could be different mechanism behind melting  $J/\psi$  and  $Y$



- Be aware of the meaning of  $T_c$  !

In quenched QCD  $T_c = 270 \text{ MeV}$

Boyd et al 1996

In full QCD there is  $T_{\text{chiral}} \approx 157 \text{ MeV}$  and  $T_{\text{onset of screening}} \approx 190 \text{ MeV}$

Budapest-Wuppertal 2010, HotQCD 2011

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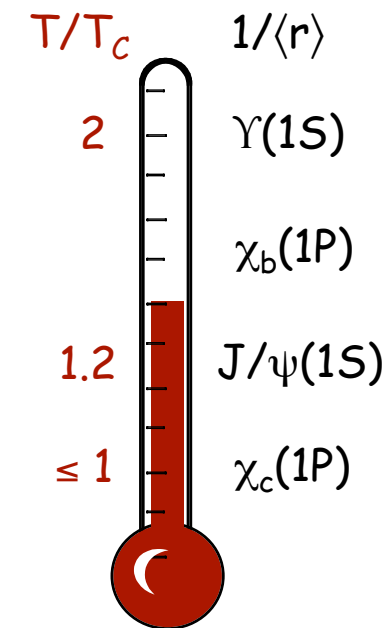
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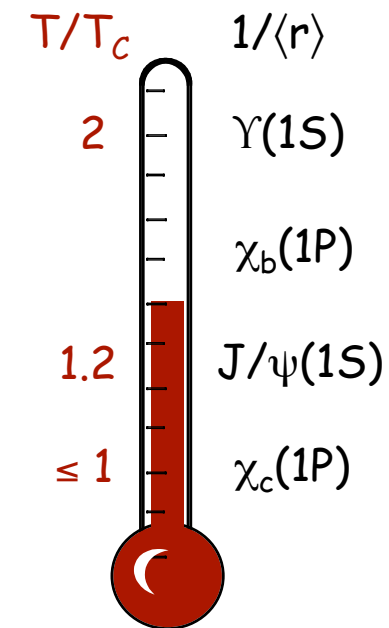
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Bazavov, Petreczky, 2010

- While there is a hierarchy in the “melting” temperatures this does not clearly imply a sequential suppression pattern

Ex: a correlated  $b$ - $\bar{b}$  pair can become a 1S or a 2S state



*“Lattice says so ...”*

- *Ugly rumor : “Lattice tells  $J/\psi$  survives to  $2T_c$ ”*

*There is no evidence for that.*

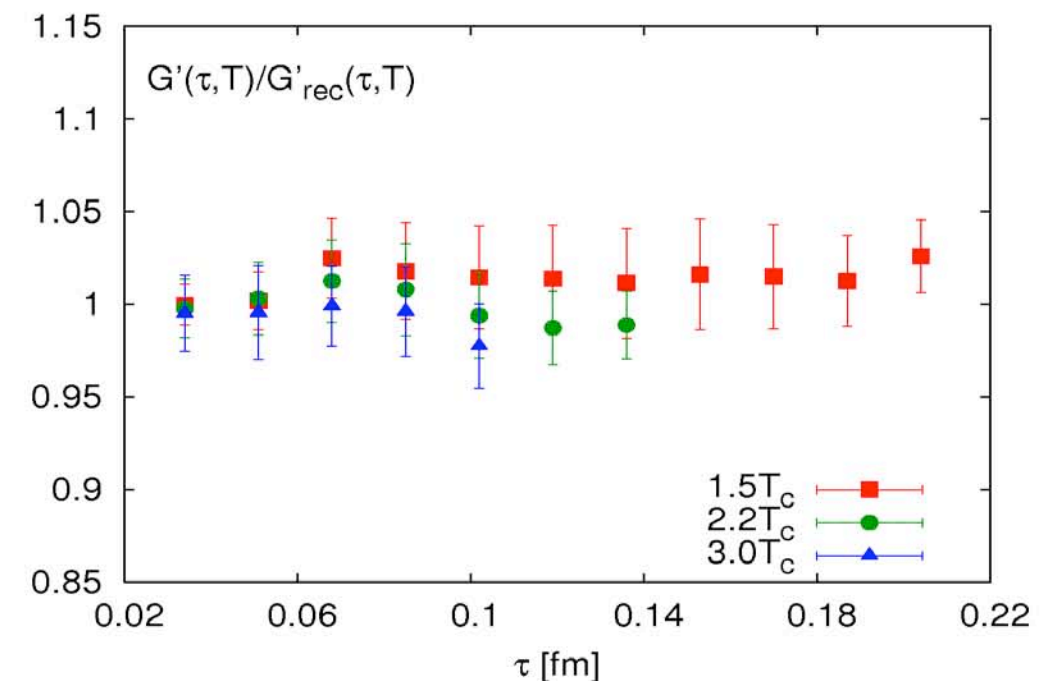
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## Correlator ratios not sensitive to spectral function changes.

- We can calculate the Euclidean correlator quantity also calculated on the lattice
- $$\frac{G(\tau, T) = \int \sigma(\omega, T) K(\tau, \omega, T) d\omega}{G_{rec}(\tau, T) = \int \sigma(\omega, T = 0) K(\tau, \omega, T) d\omega}$$
- Correlators do not change just as lattice says so!
  - We now understand that changes come from zero modes



Datta et al; Jakovac et al; Umeda; Petreczky, 2005-2009

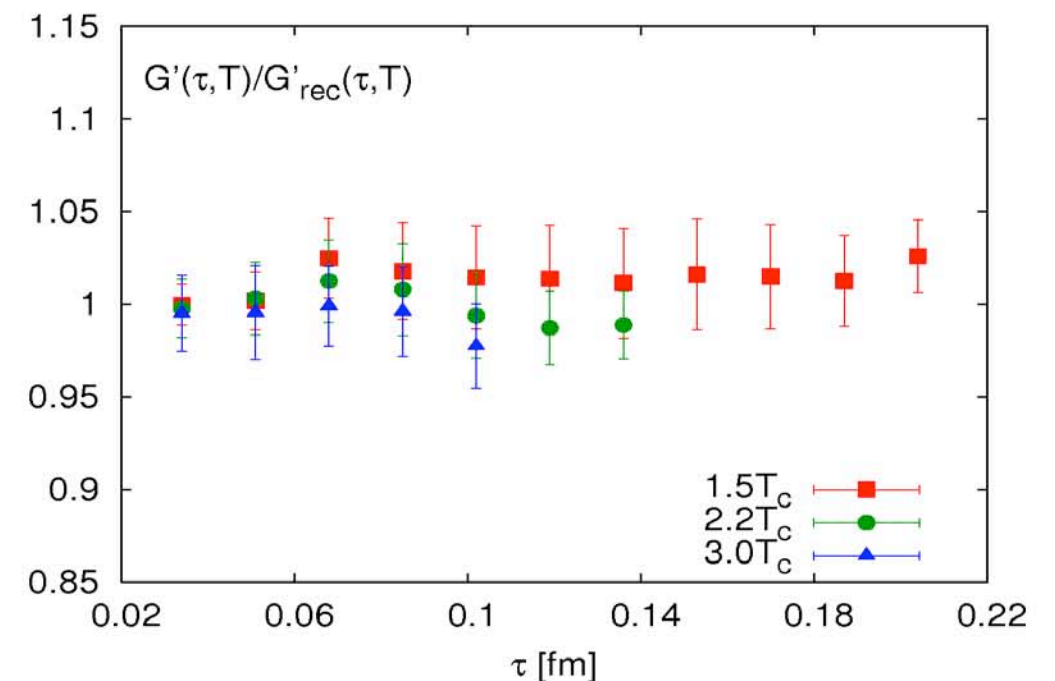
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+ Lattice spectral functions do not suggest  $J/\psi$  survival

# Charmonium Spectral Function

Extracted in quenched QCD

Extracted from correlation function of mesonic currents in Euclidean time

$$G(\tau, T) = \int \sigma(\omega, T) K(\tau, \omega, T) d\omega$$

correlator  
*directly calculated*

spectral function extracted  
*not directly calculated*

$\mathcal{O}(10)$  data but  $\mathcal{O}(100)$  degrees of freedom to reconstruct



maximizes the conditional probability  
of having the spf given the data and  
some prior knowledge

Umeda et al, EPJ C39S1 (05) 9, Asakawa,  
Hatsuda, PRL 92 (2004) 01200, Datta et al, PRD  
69 (04) 094507, Jakovac et al PRD 2007, ...

Shortcomings: limited # of data points

limited extent in tau

default model dependence *large*

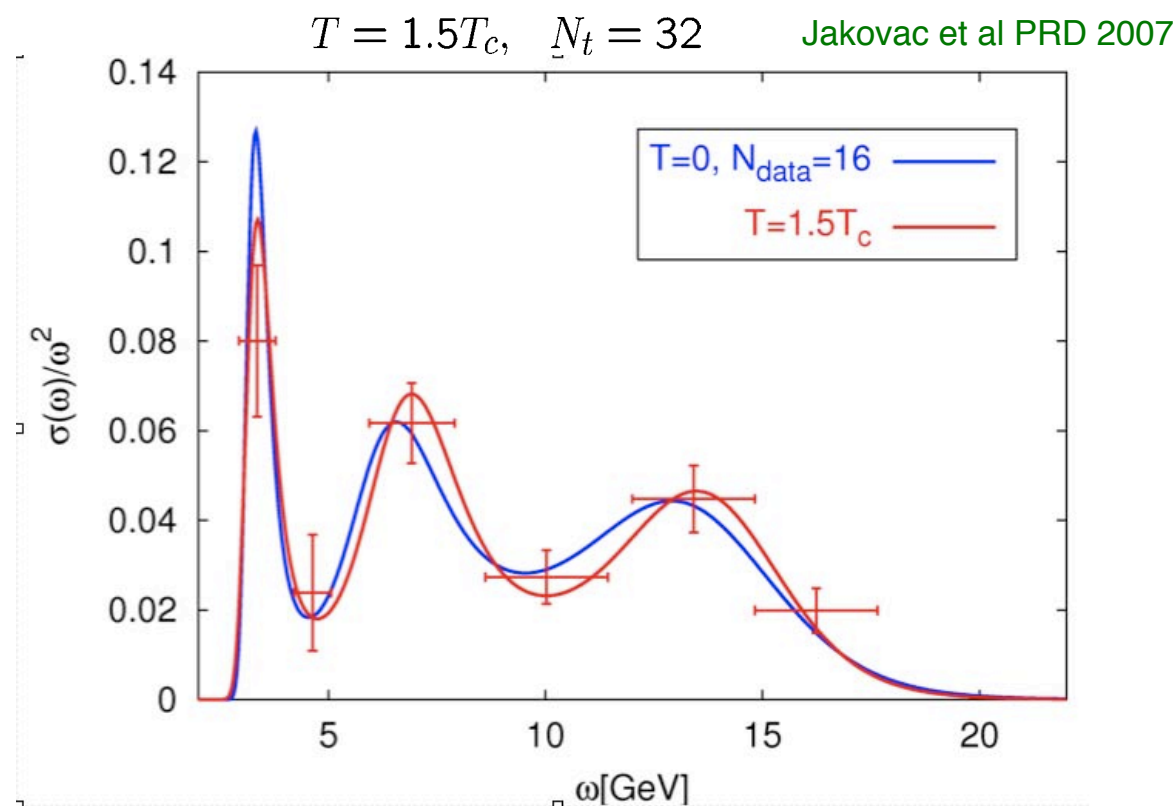
$$T = 1/(N_t a)$$

$$\tau_{max} = 1/(2T)$$



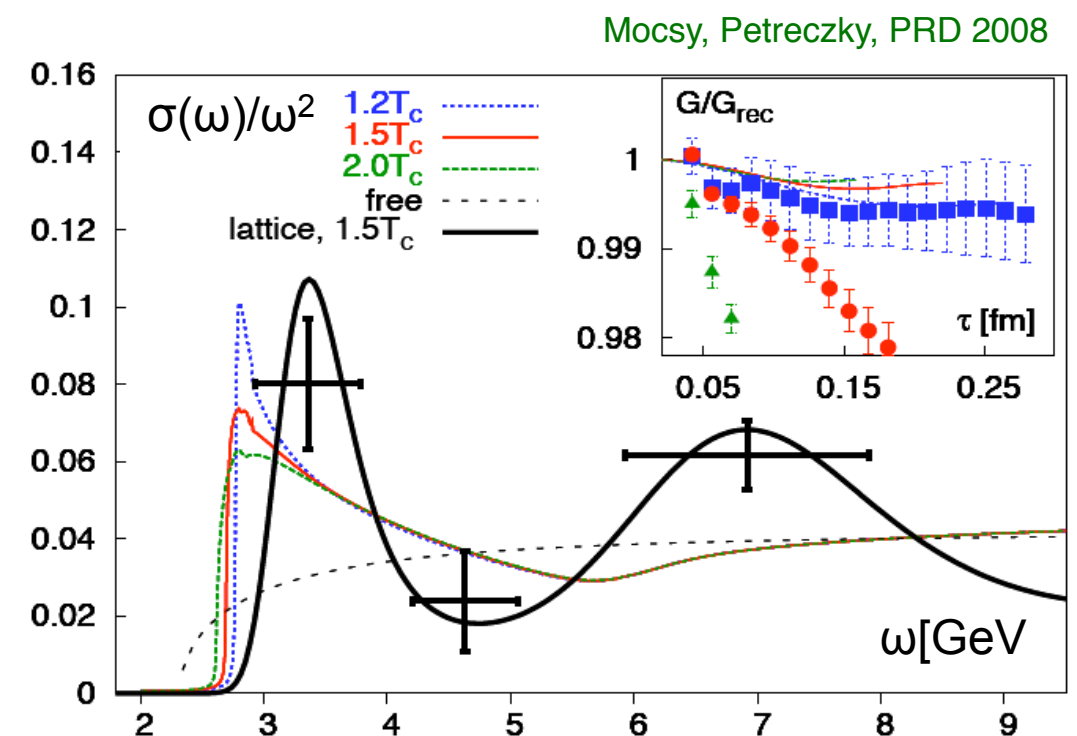
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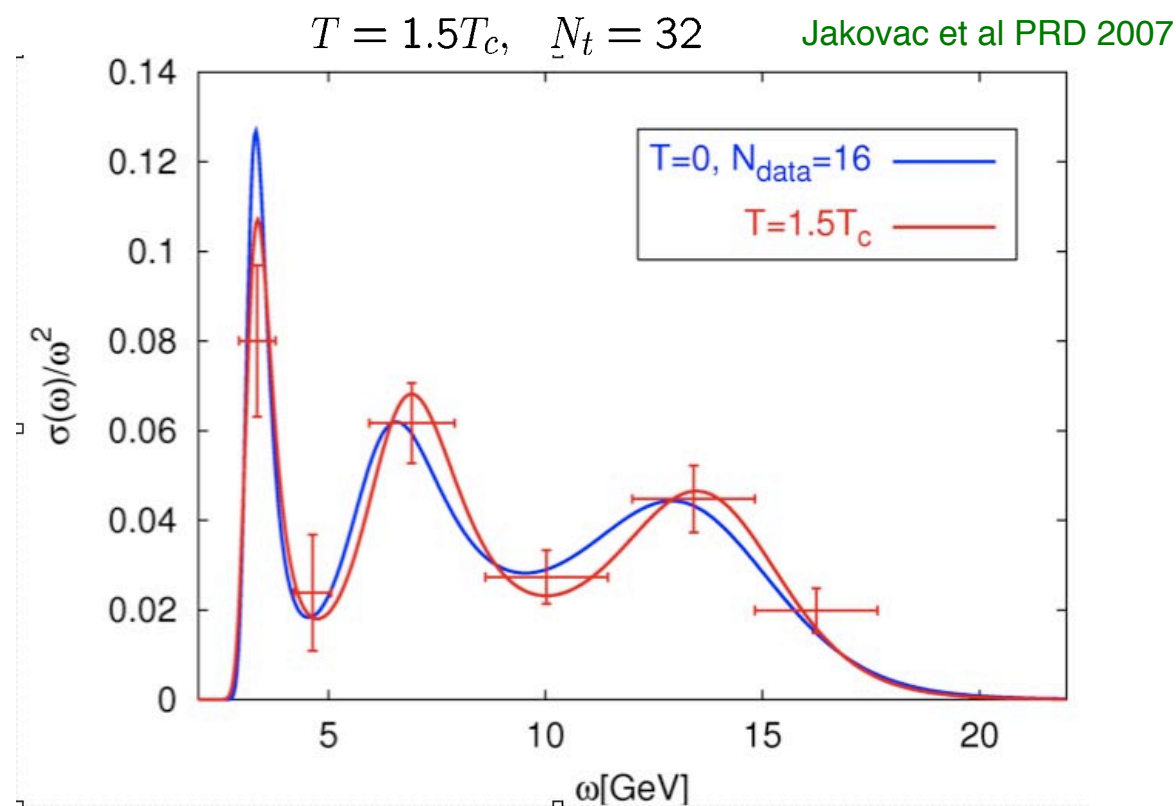
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1st peak consistent with threshold enhancement

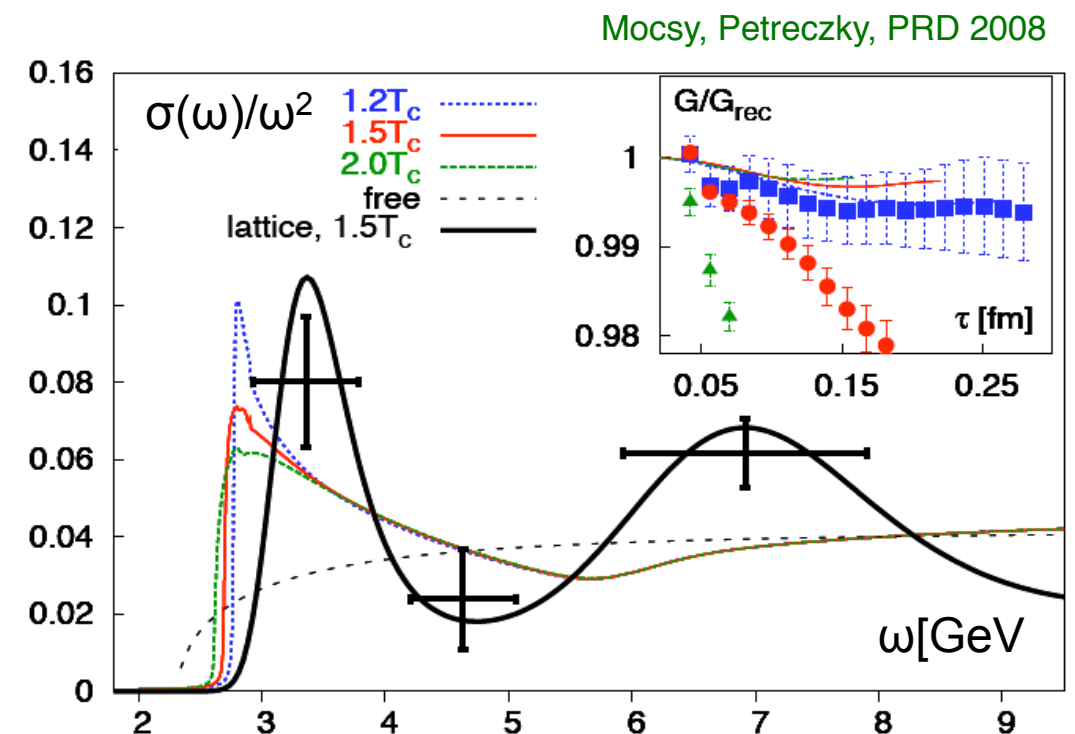
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Comparing low resolution confined phase (blue) to low resolution deconfined phase (red) and getting an agreement does **not** imply the agreement will hold at high resolution

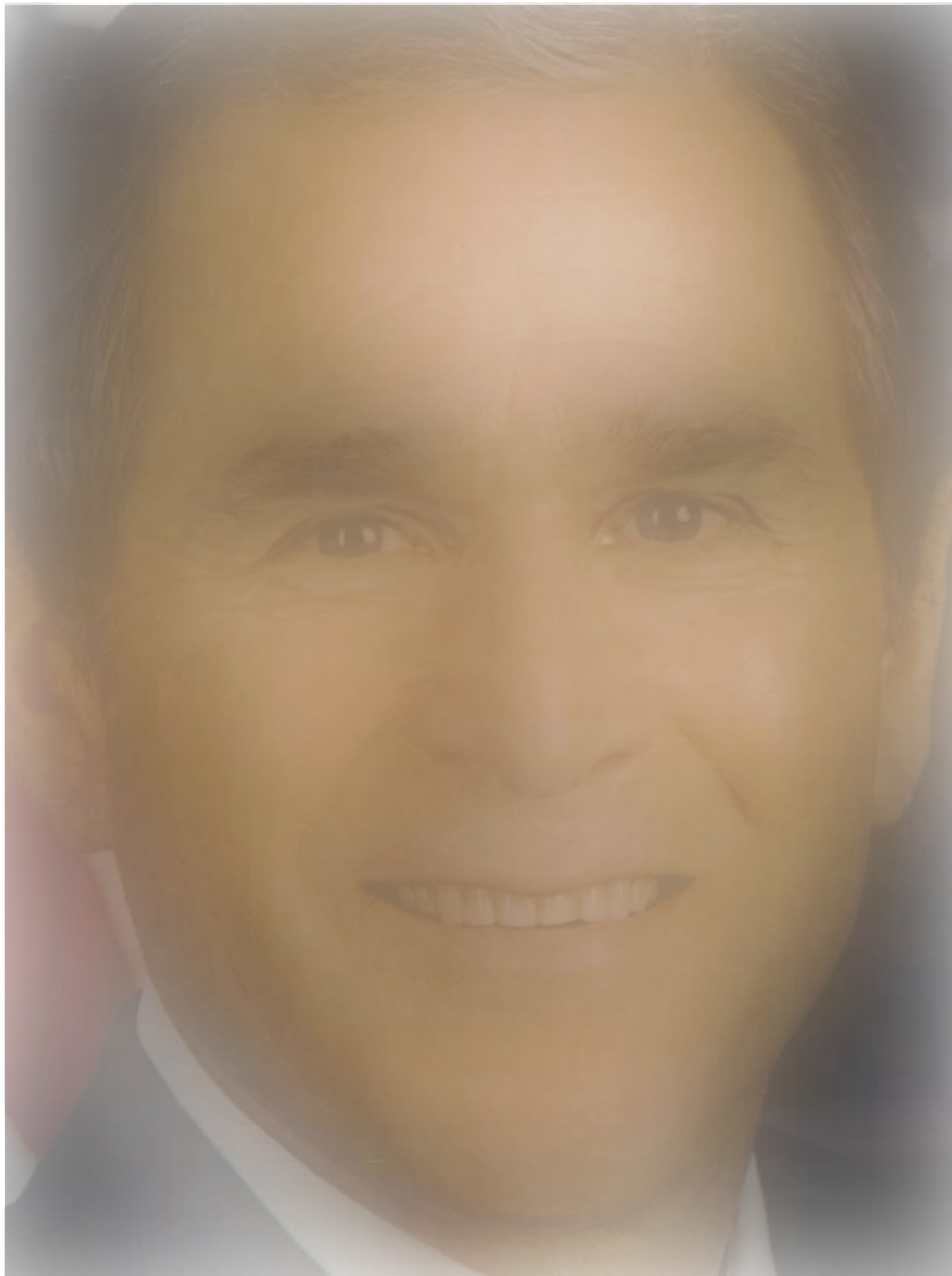


agreement in low resolution does  
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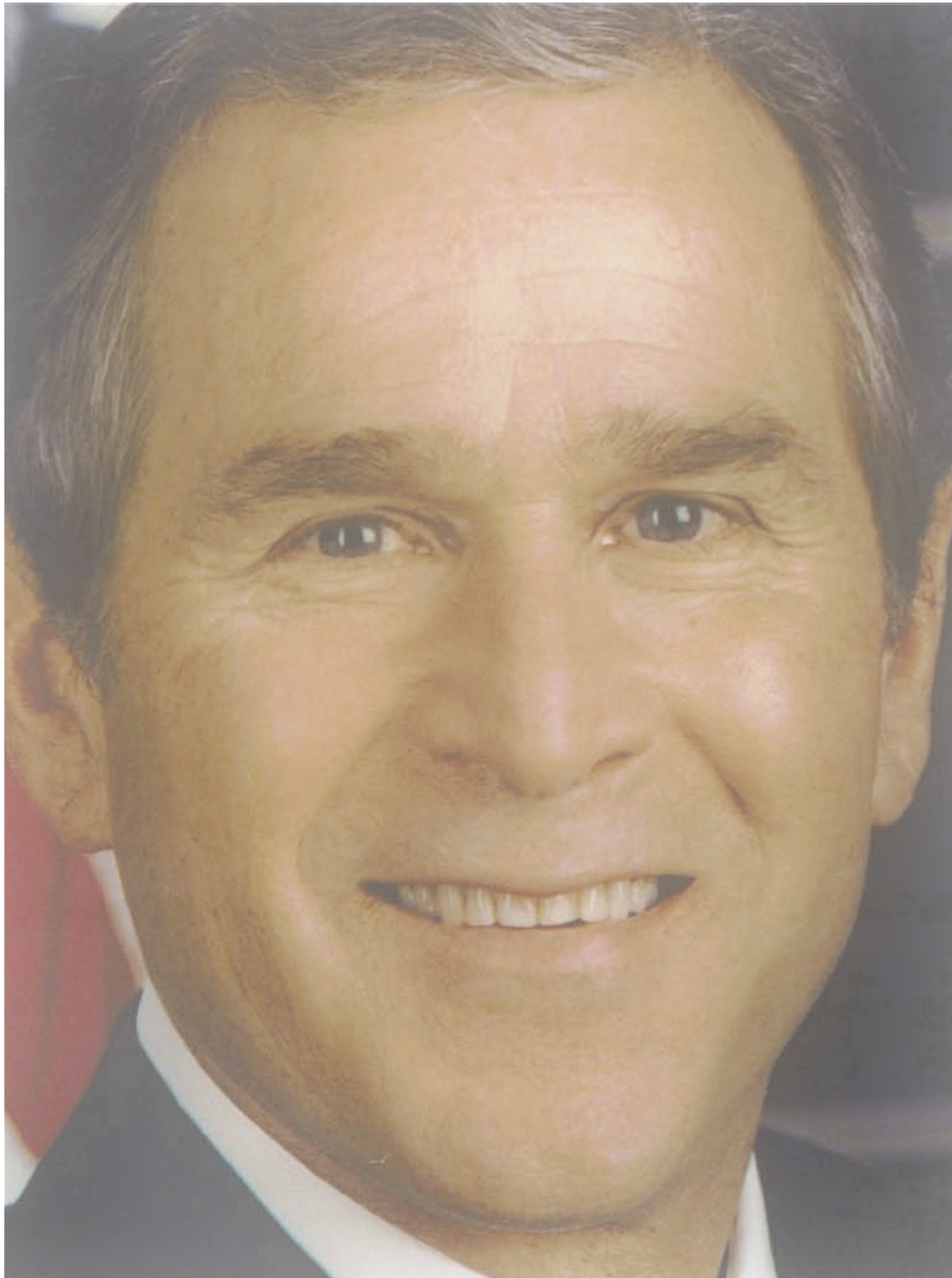


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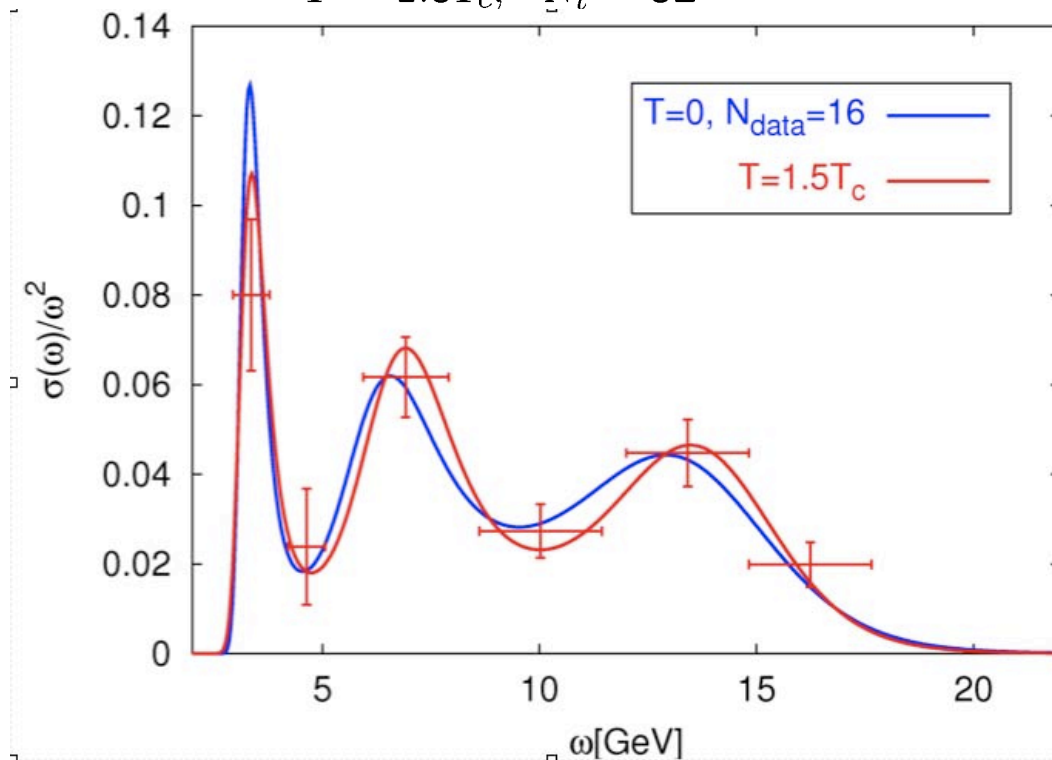
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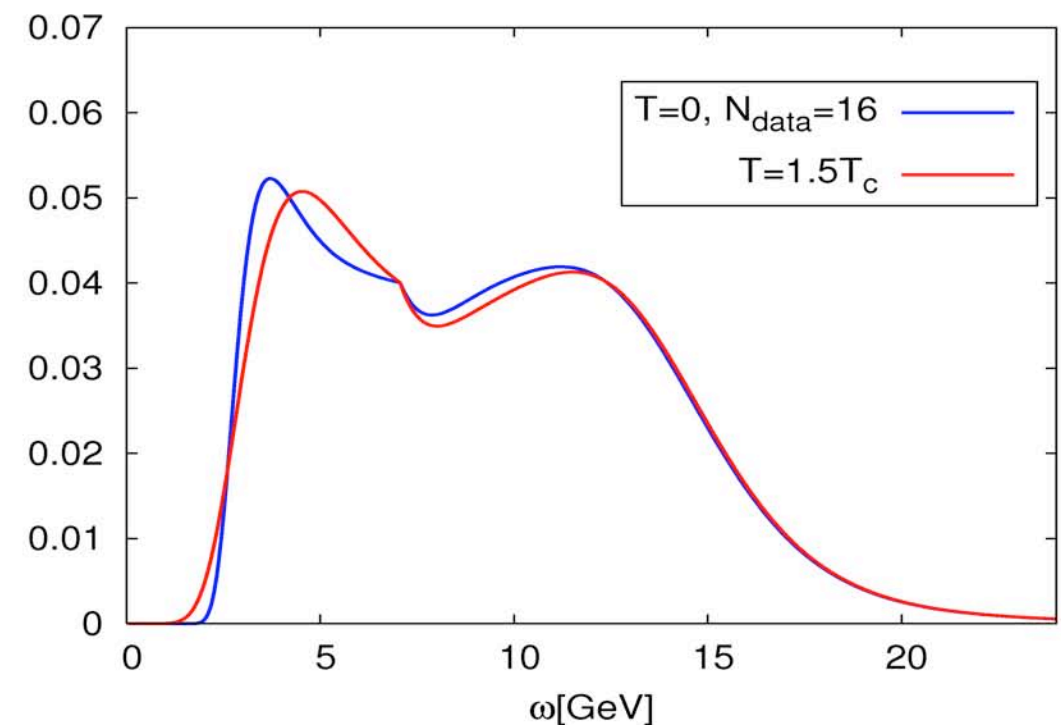
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$T = 1.5T_c$ ,  $N_t = 32$  Jakovac et al PRD 2007



prior: perturbative continuum spf (no lattice effects)



prior: from  $T=0$  data at high energies

$$T = 1/(N_t a)$$

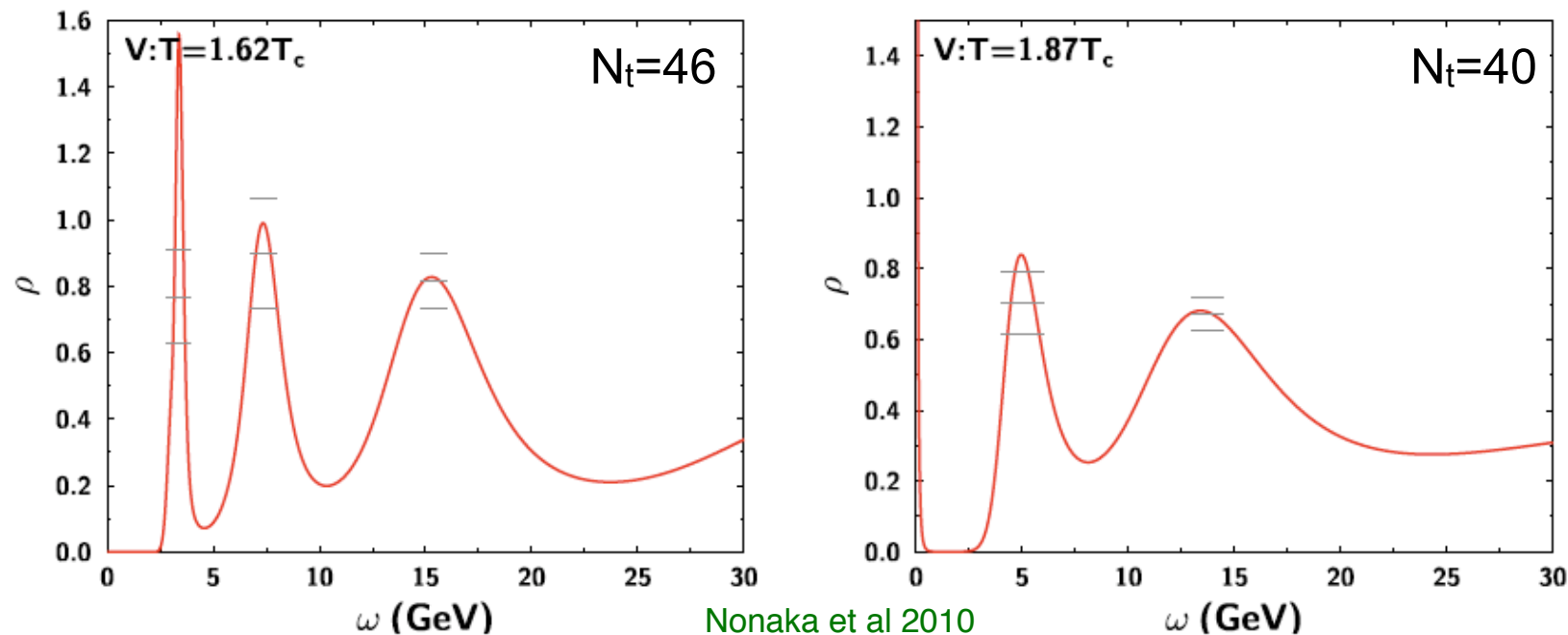
$$\tau_{max} = 1/(2T)$$

Strong default model dependence  
peak - no peak ?!

# Charmonium Spectral Function

Extracted in quenched QCD

## Most recent lattice results



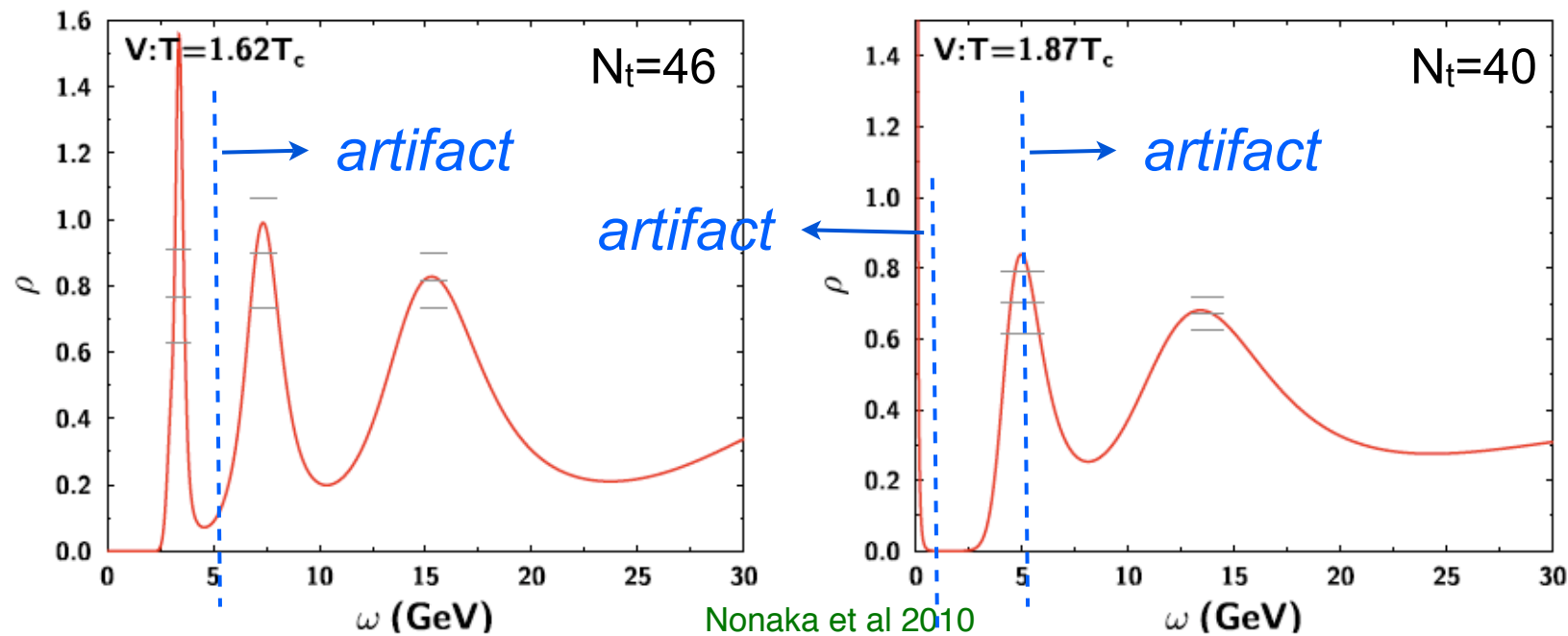
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“We found that  $J/\psi$  is stable above the critical temperature  $T=1.62T_c$ ”

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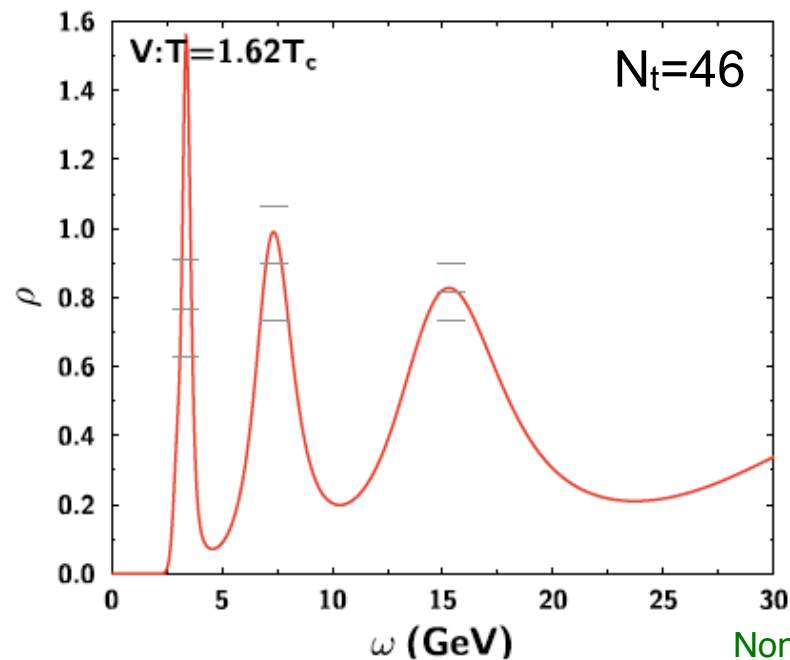
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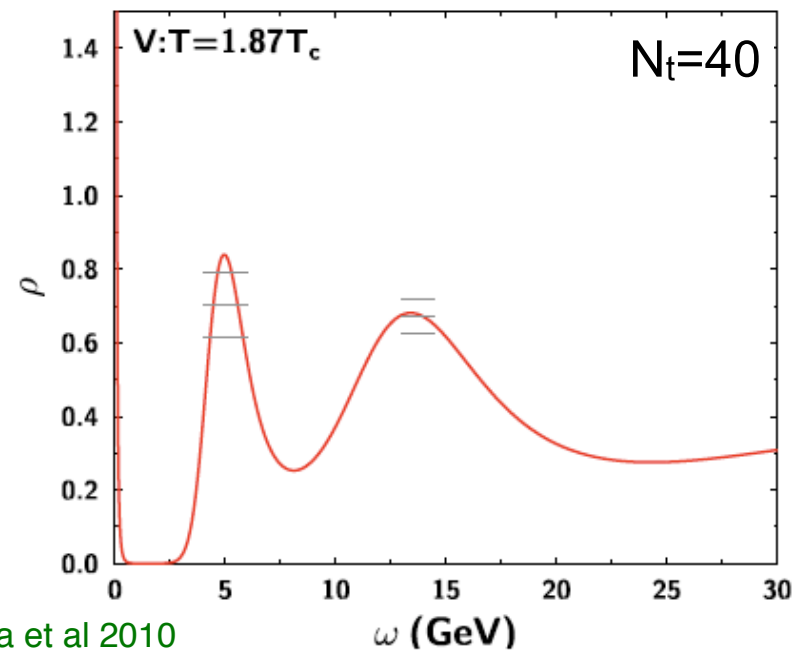
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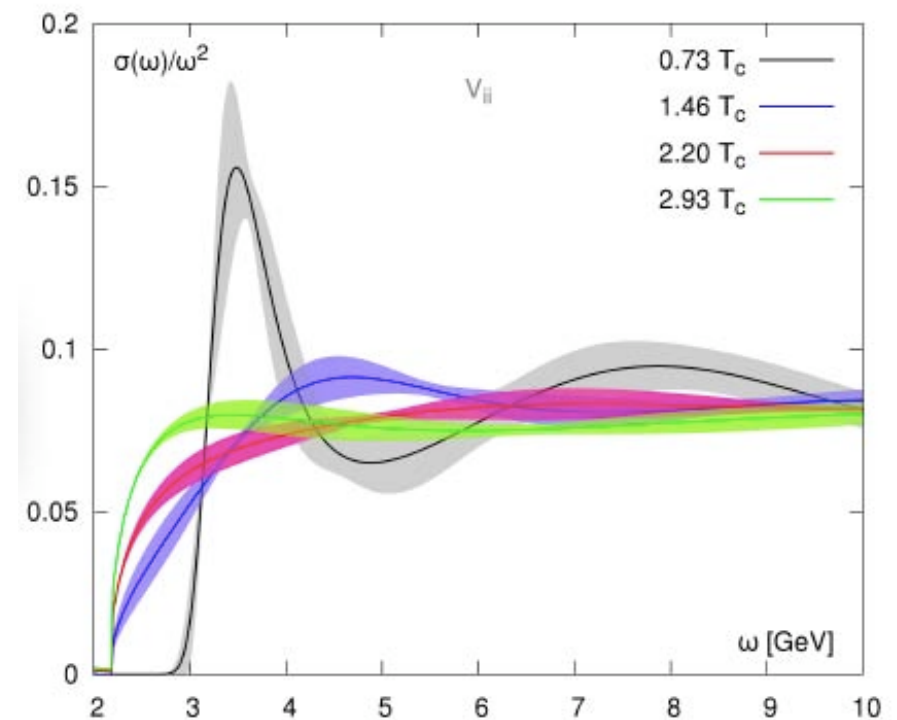


Nonaka et al 2010



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Ding et al 2010

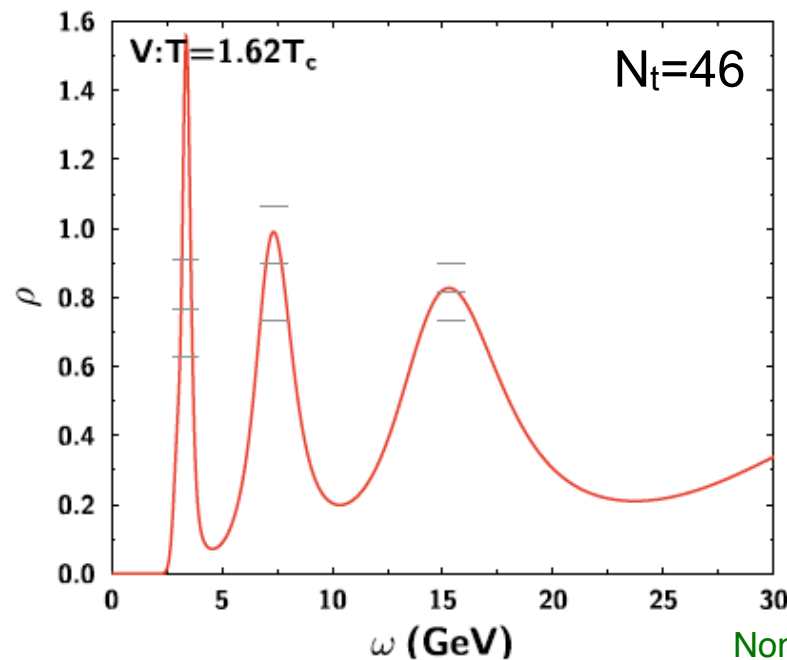
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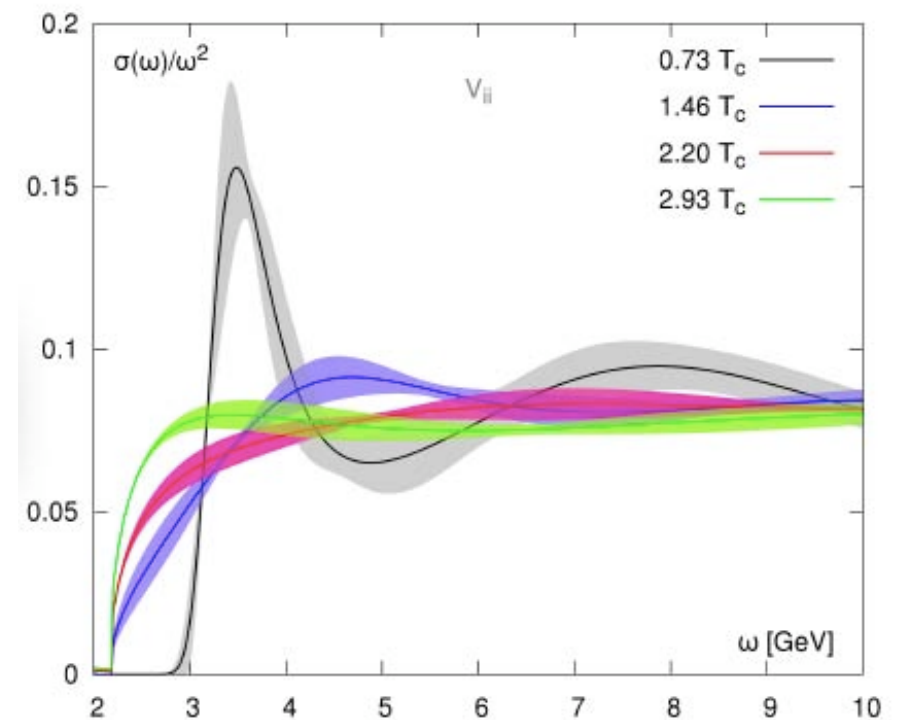
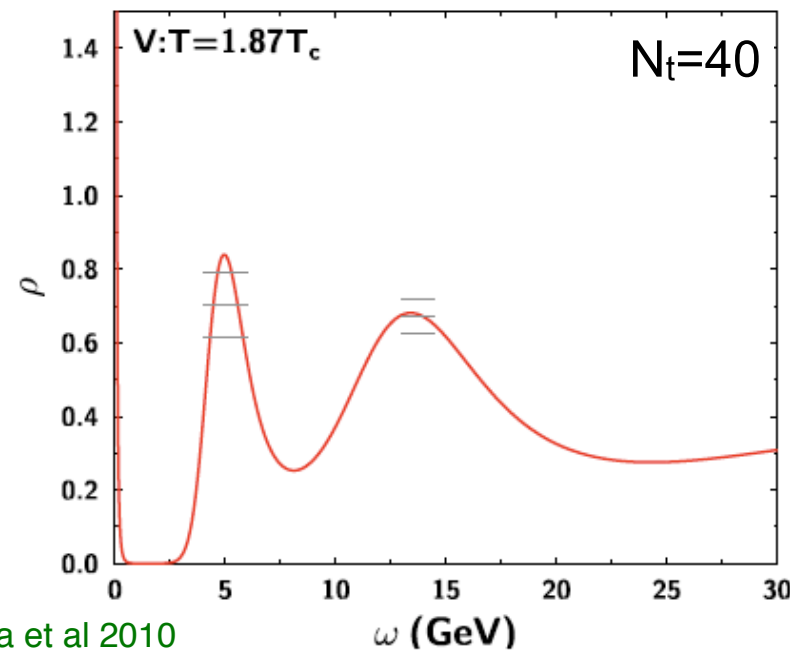
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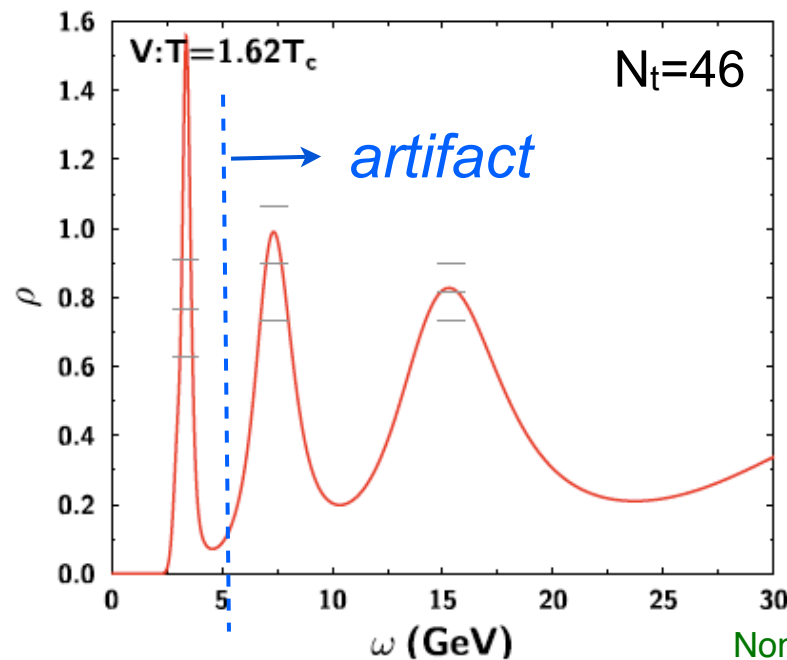
“Our analysis suggests that  $J/\psi$  is melted already by  $1.46 T_c$ ”

Even if we take all the 1st peaks seriously the claims are contradictory.

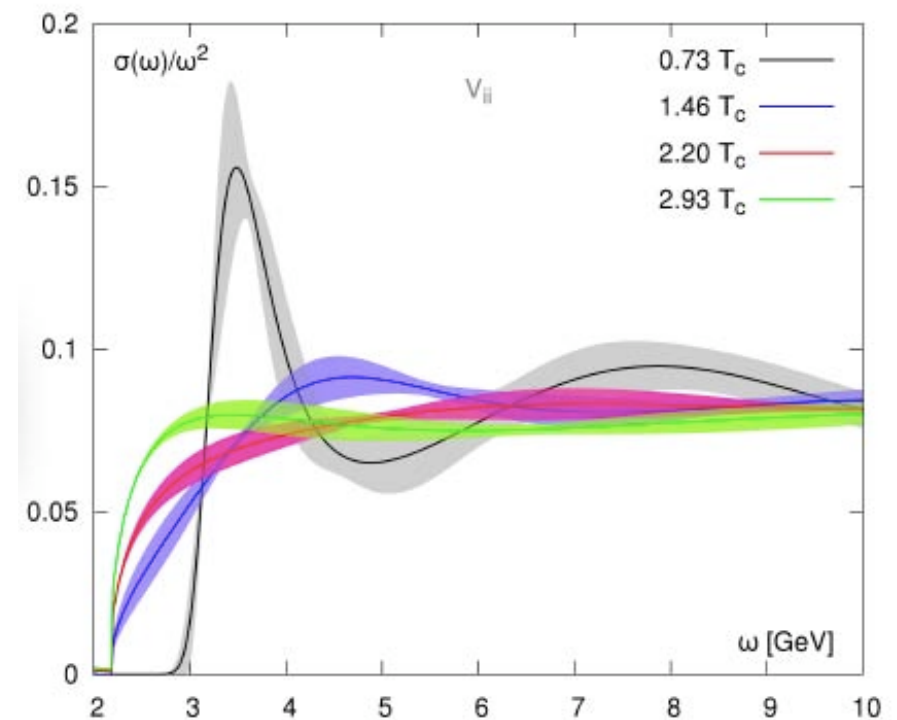
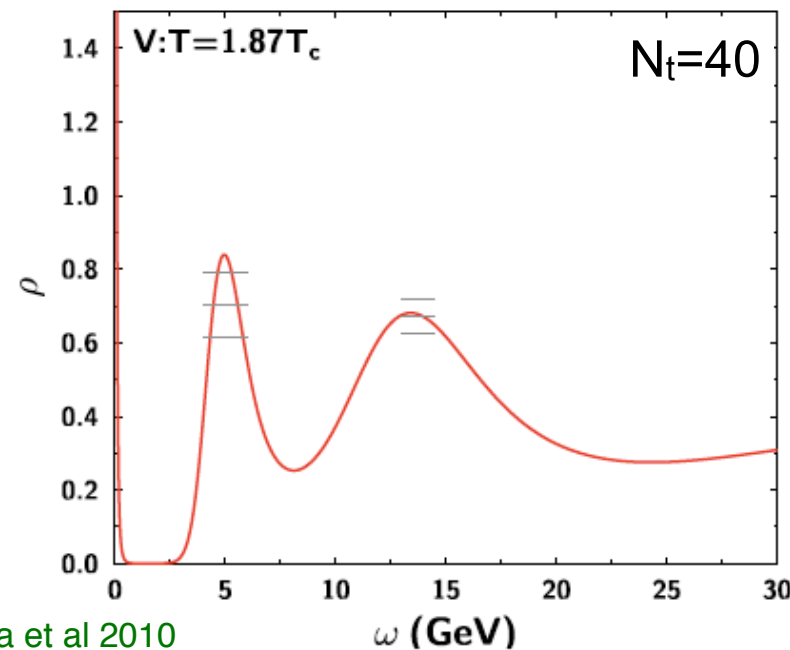
# Charmonium Spectral Function

Extracted in quenched QCD

## Most recent lattice results



Nonaka et al 2010



Ding et al 2010

prior: perturbative continuum spf (no lattice effects)

prior: from  $T=0$  data at high energies

“We found that  $J/\psi$  is stable above the critical temperature  $T=1.62T_c$ ”

“Our analysis suggests that  $J/\psi$  is melted already by  $1.46 T_c$ ”

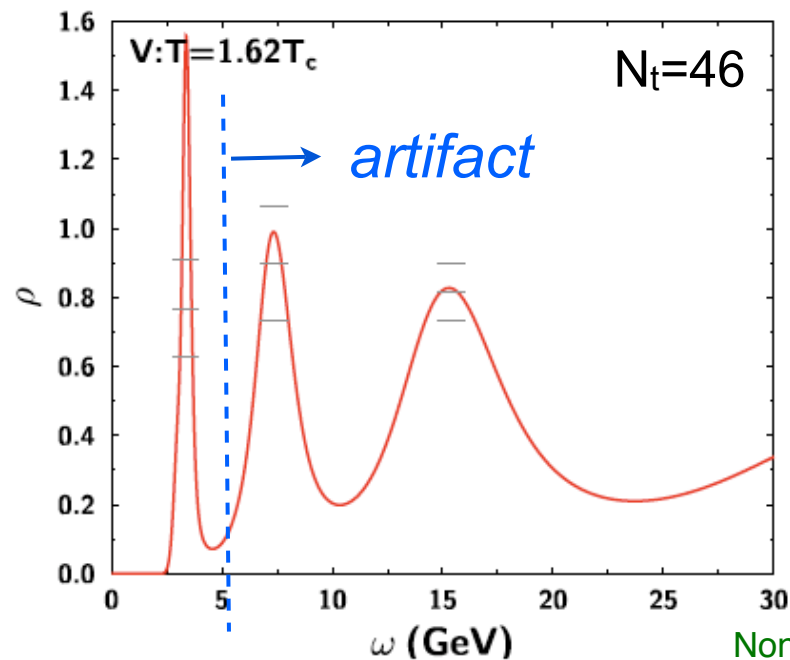
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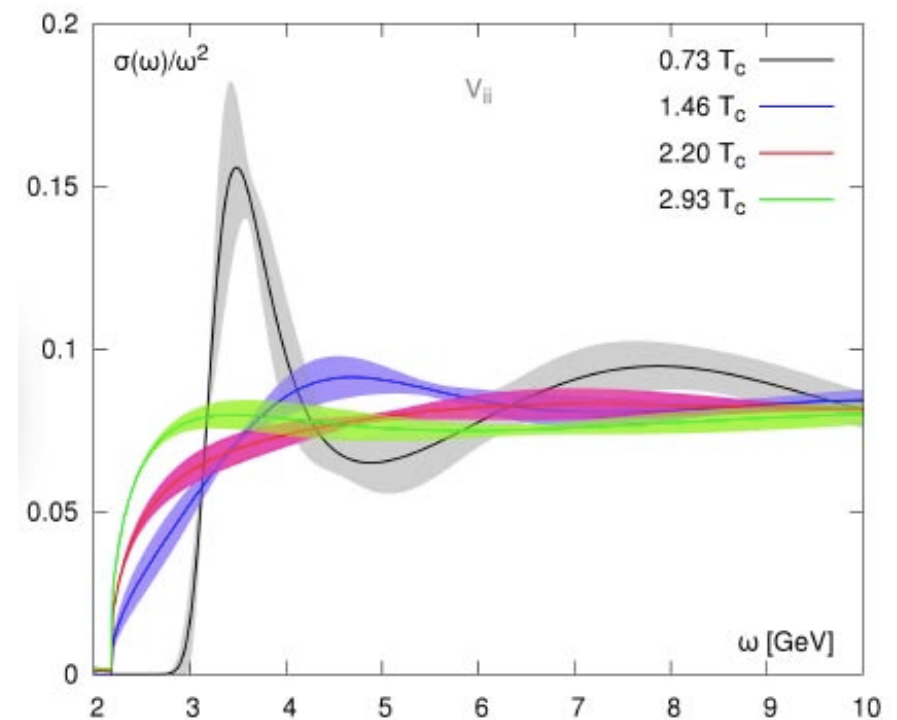
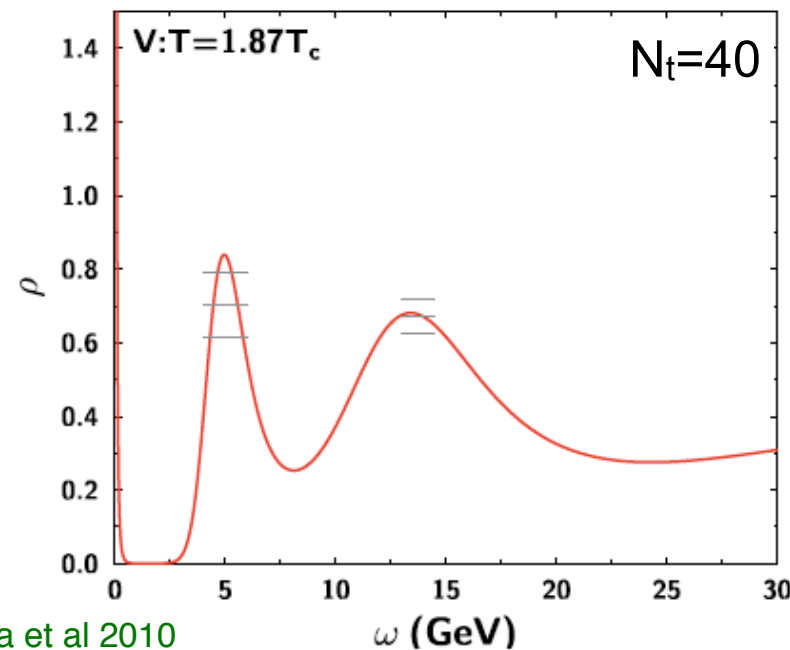
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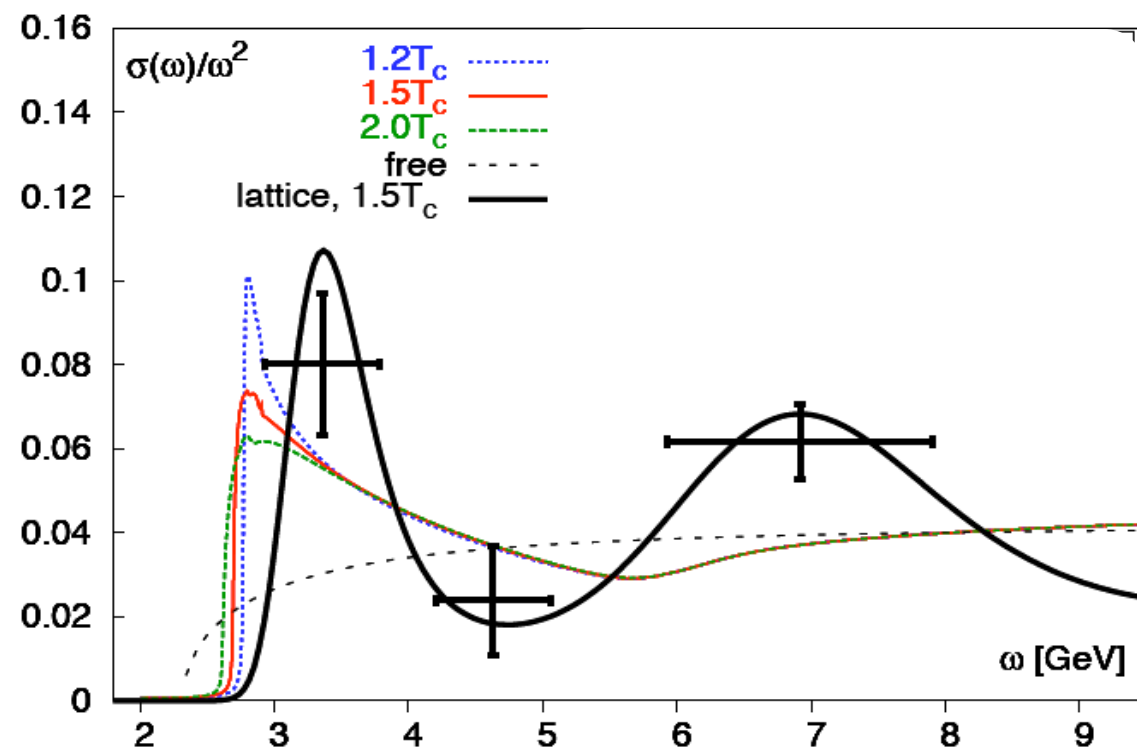
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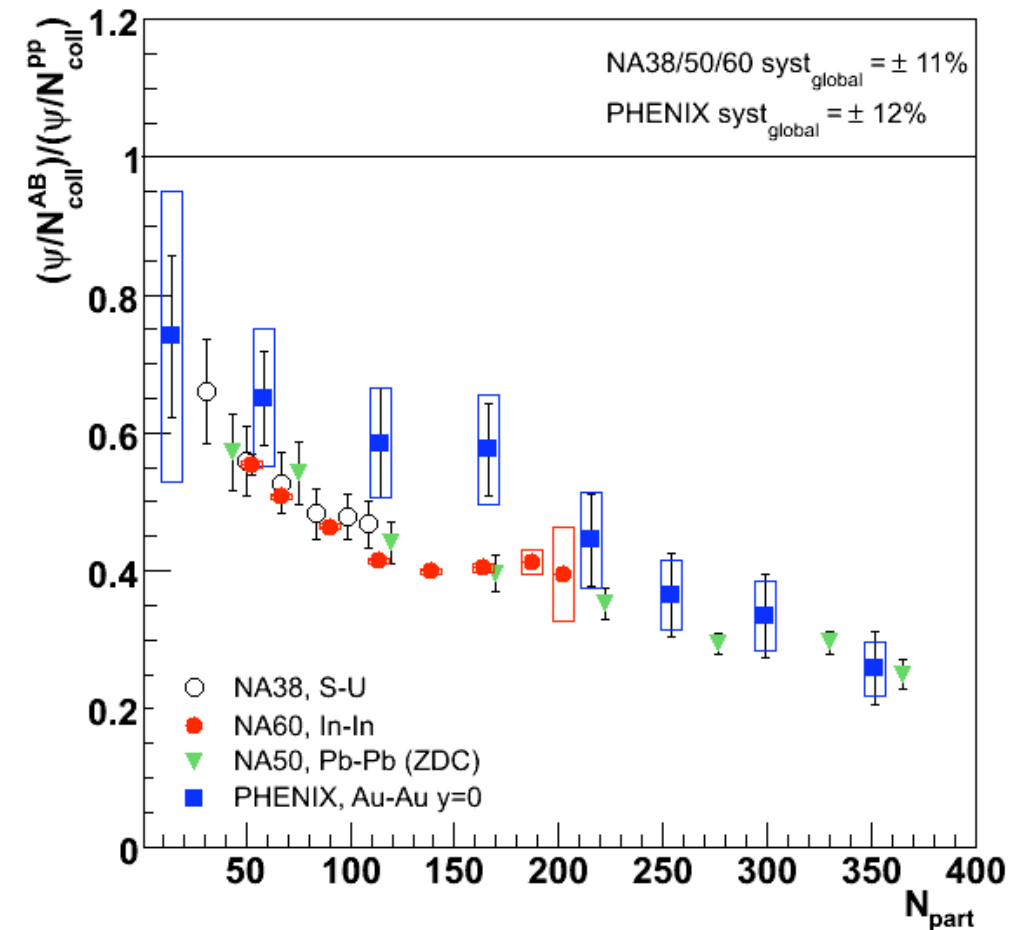
It is difficult to make any conclusive statement regarding quarkonium at finite temperature based on the shape of the lattice spectral functions

# Famous Plot

How can we relate these?



Quarkonium spectral functions  
in equilibrated plasma

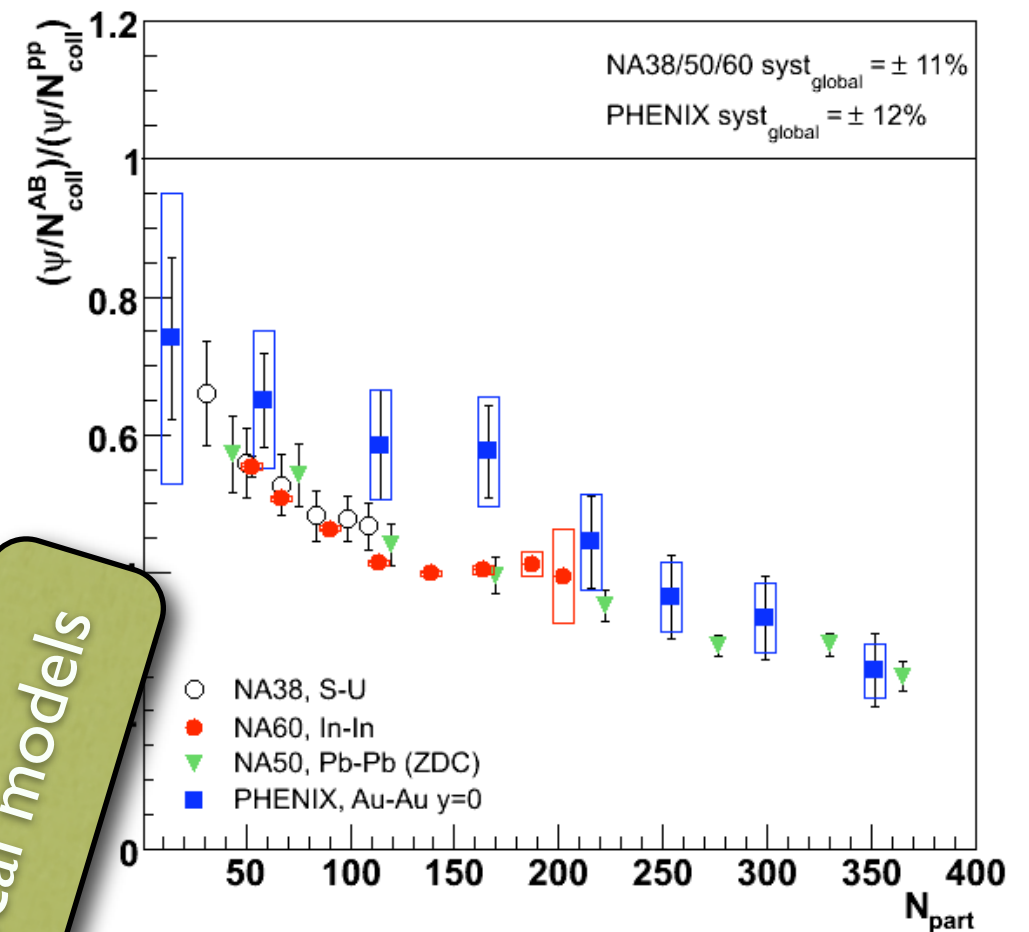
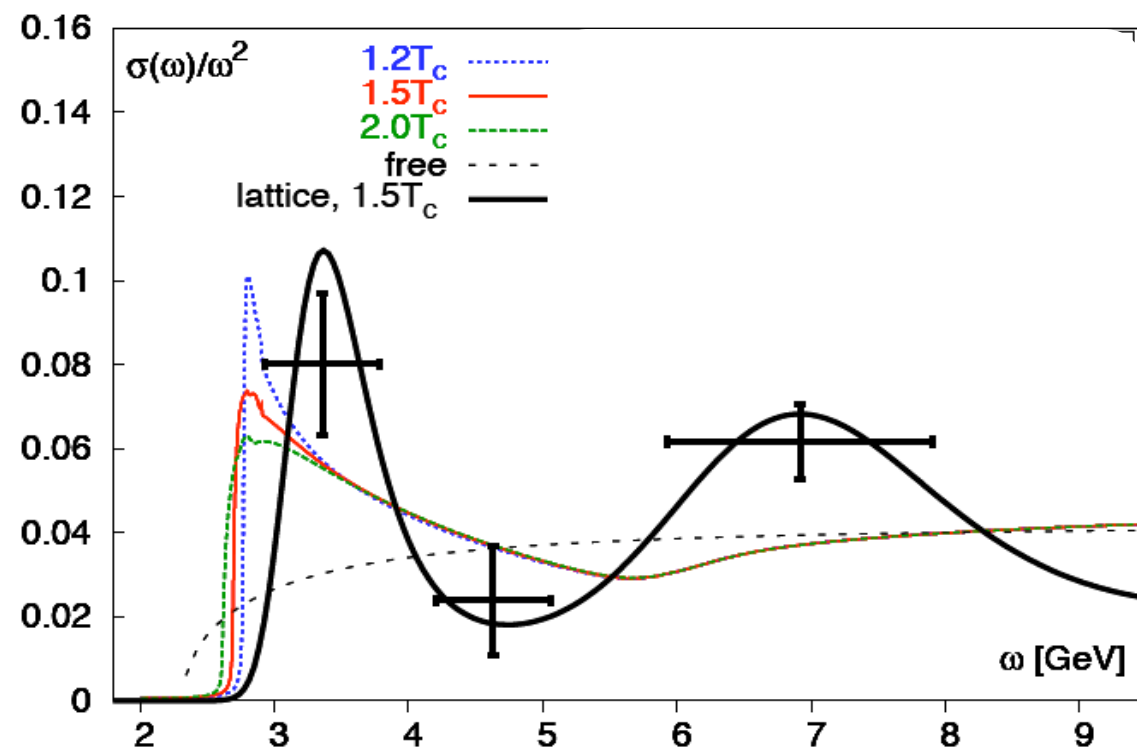


Experimental data



# Famous Plot

How can we relate these?



Quarkonium spectral functions  
in equilibrated plasma

Dynamical models

Experimental data

# The Connection

Zhao, Rapp, van Hees, Fries, Young, Shuryak, ...

- Spectral function calculation → no bound states  
only correlated c-cbar pairs
- What is the probability that c-cbar  
find themselves in proximity at the  
hadronization time?
- Modeling the motion of c-cbar in the  
evolving fireball  
according to Langevin dynamics
  - stochastic force from the heat bath
    - Input: charm diffusion constant  
Best current guess Teaney and Moore  $D_c(2\pi T) = 1.5-3$   
if small enough that attraction between c and cbar may survive
  - attractive interaction between c-cbar
    - Input: heavy quark potential
  - lifetime of the plasma

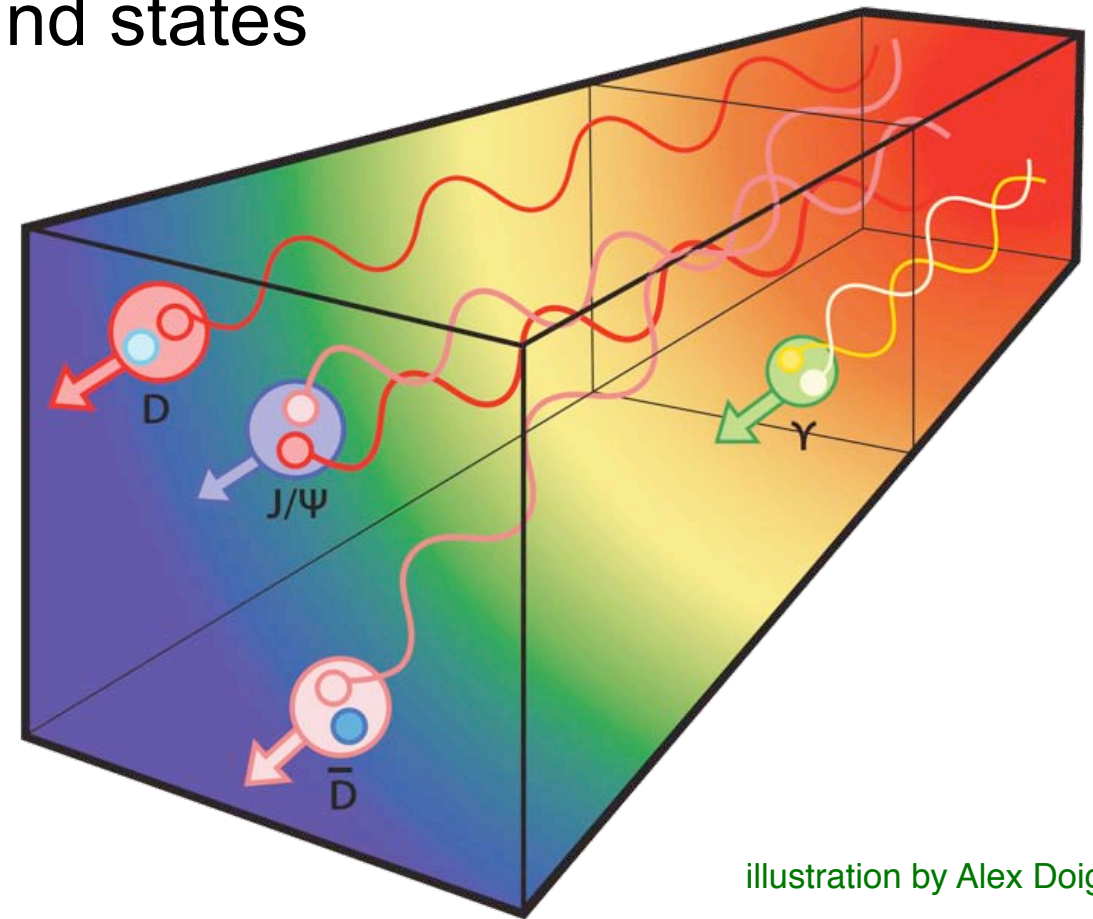
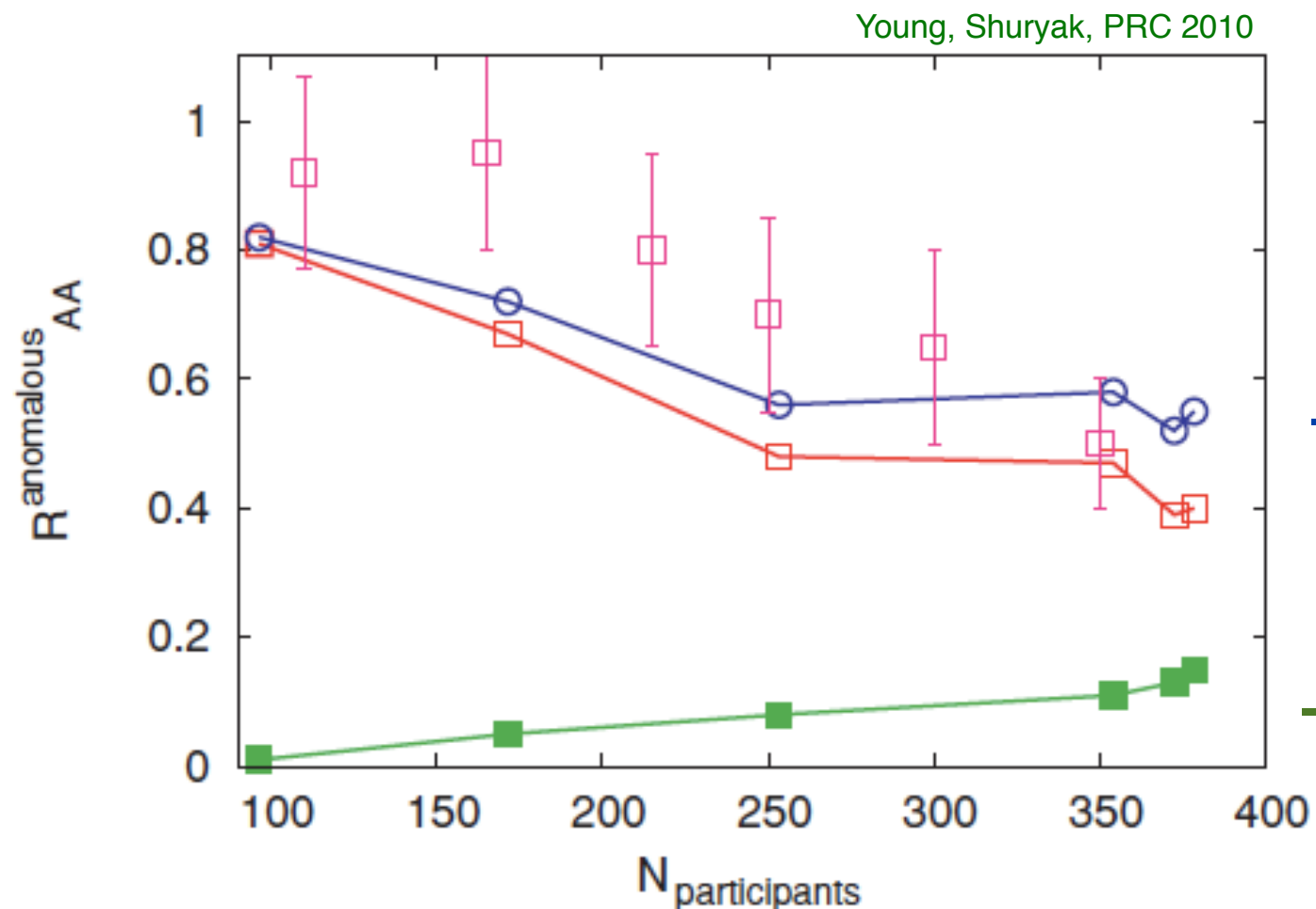


illustration by Alex Doig

# The Connection

## - Comparison to PHENIX data



- total yield

- suppression +  
correlated regeneration

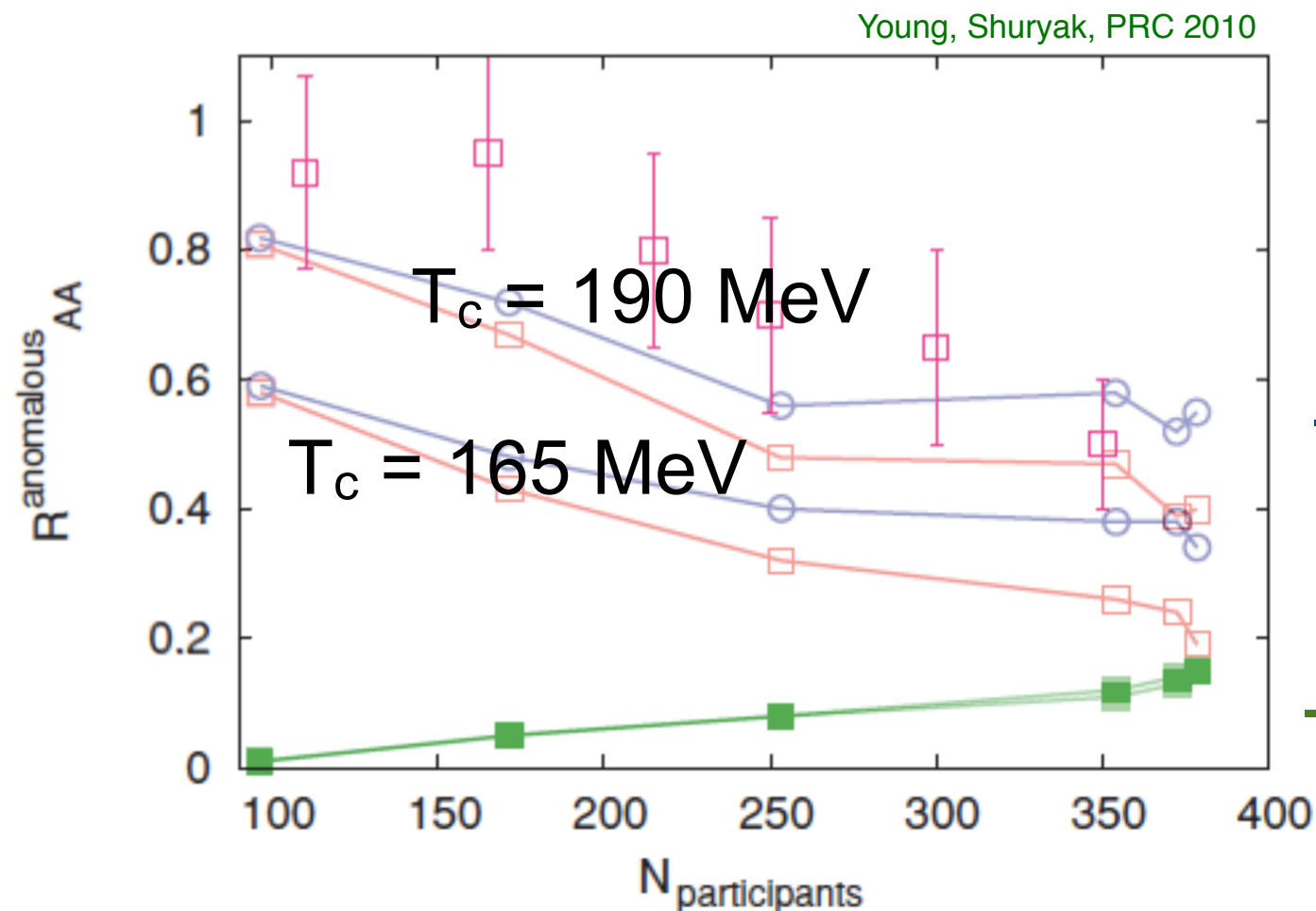
- statistical recombination  
c and cbar originate from  
different hard processes

- Direct  $J/\psi$  suppressed, but  $\sim 50\%$  of correlated c-cbar recombine
- Coalescence gives relative small contribution
- Quite good agreement with data for small charm diffusion and  $T_c=190$  MeV

Note:  $T_c = T_{\text{deconfinement}}$

# The Connection

## - Comparison to PHENIX data



- total yield

- suppression +  
correlated regeneration

- statistical recombination  
c and cbar originate from  
different hard processes

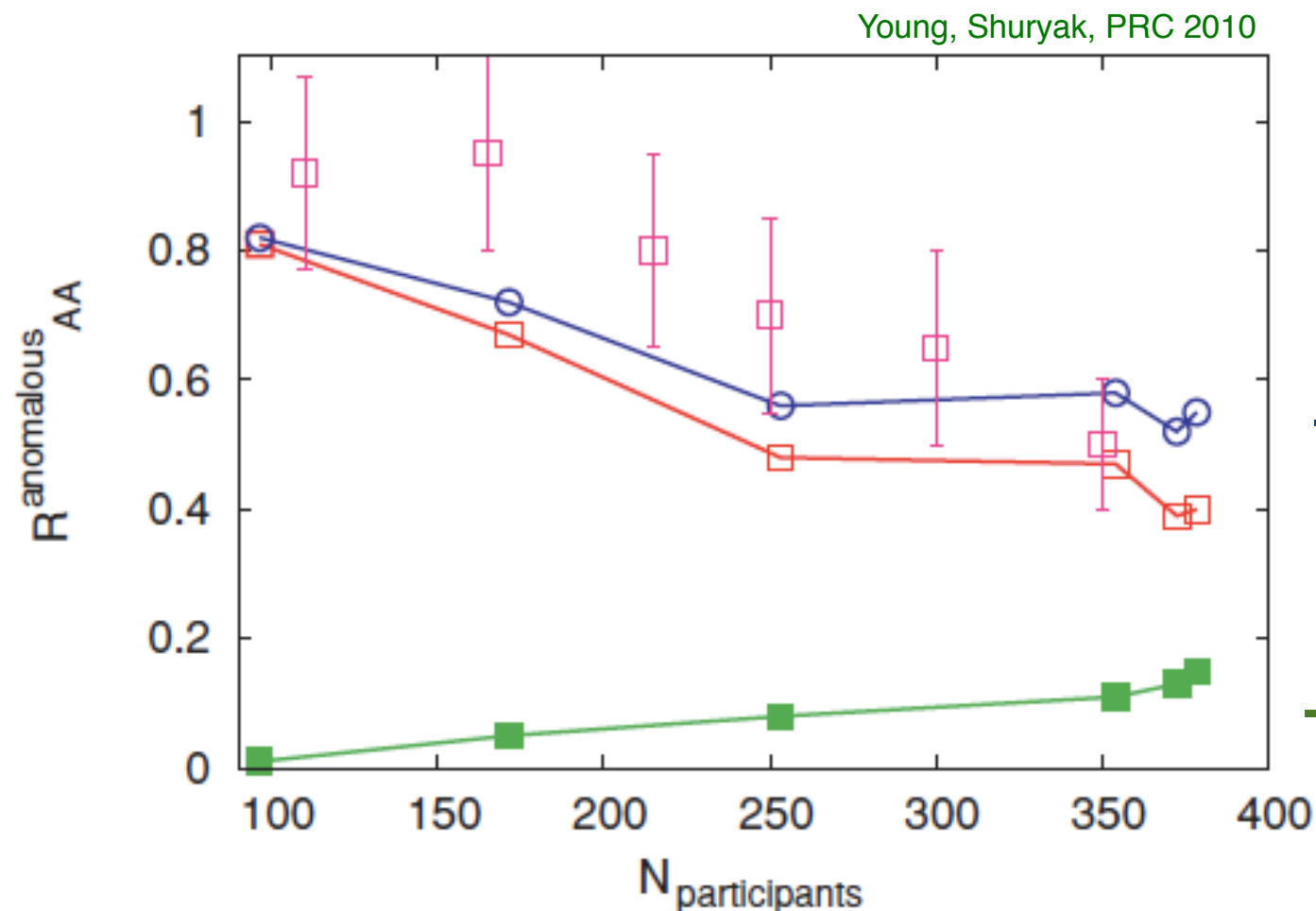
- Agreement is better with the higher deconfinement temperature,  
i.e. with a shorter QGP lifetime

The finiteness of  $R_{AA}$  can lead us to determine the plasma lifetime ?!

Note:  $T_c = T_{\text{deconfinement}}$

# The Connection

- Comparison to PHENIX data



- total yield

- suppression + correlated regeneration

- statistical recombination  
c and cbar originate from different hard processes

- Note

There are effects not included in this model: initial state effects and absorption in the crossover-hadronic region (CNM effects)

- A quantitative comparison with data is difficult



# Can we isolate hot effects?

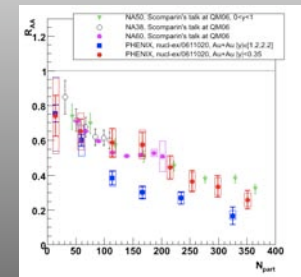
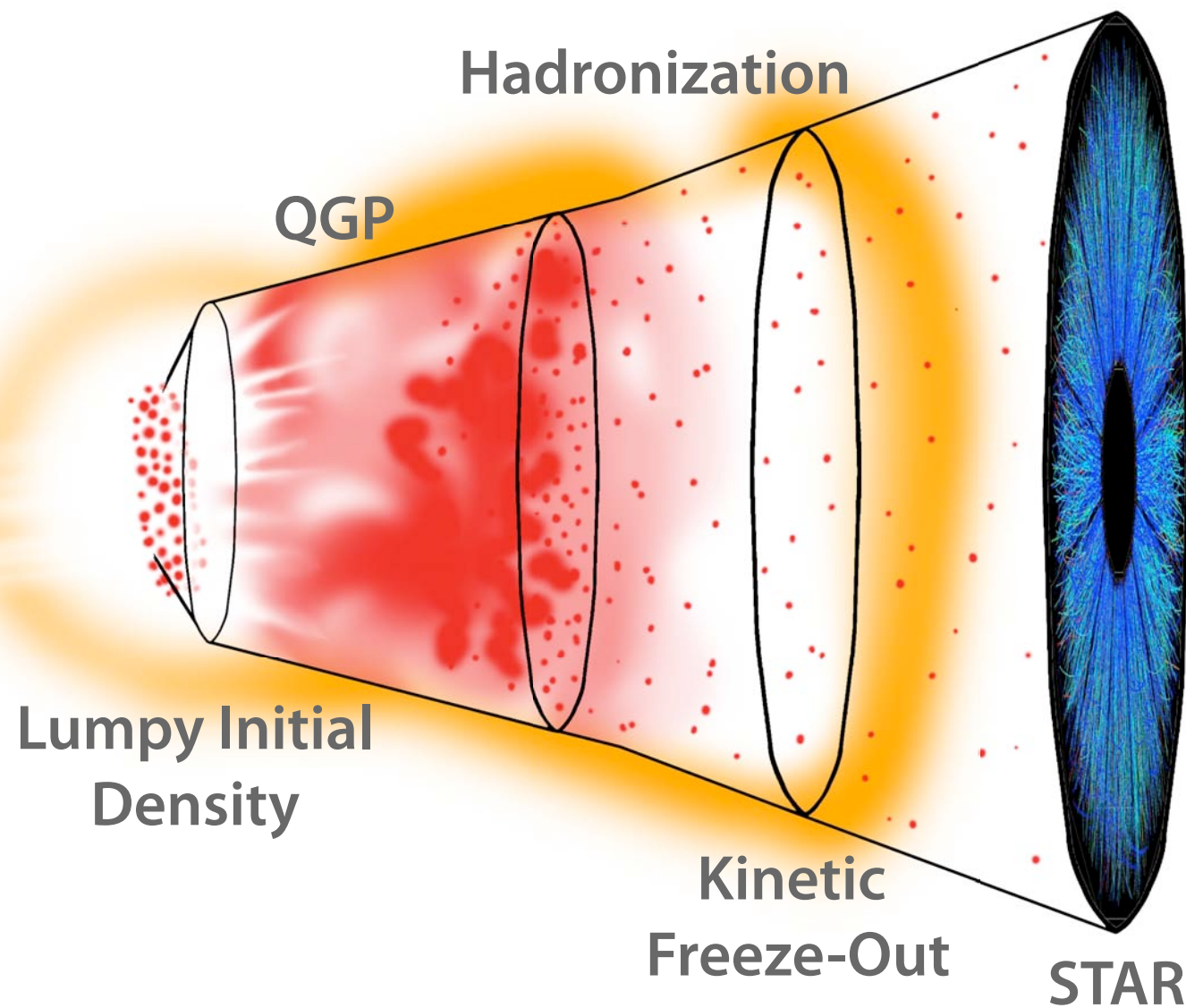


illustration by Alex Doig

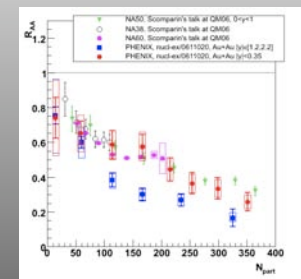


c and cbar are  
produced at  
early times ...

... go through the entire evolution ...

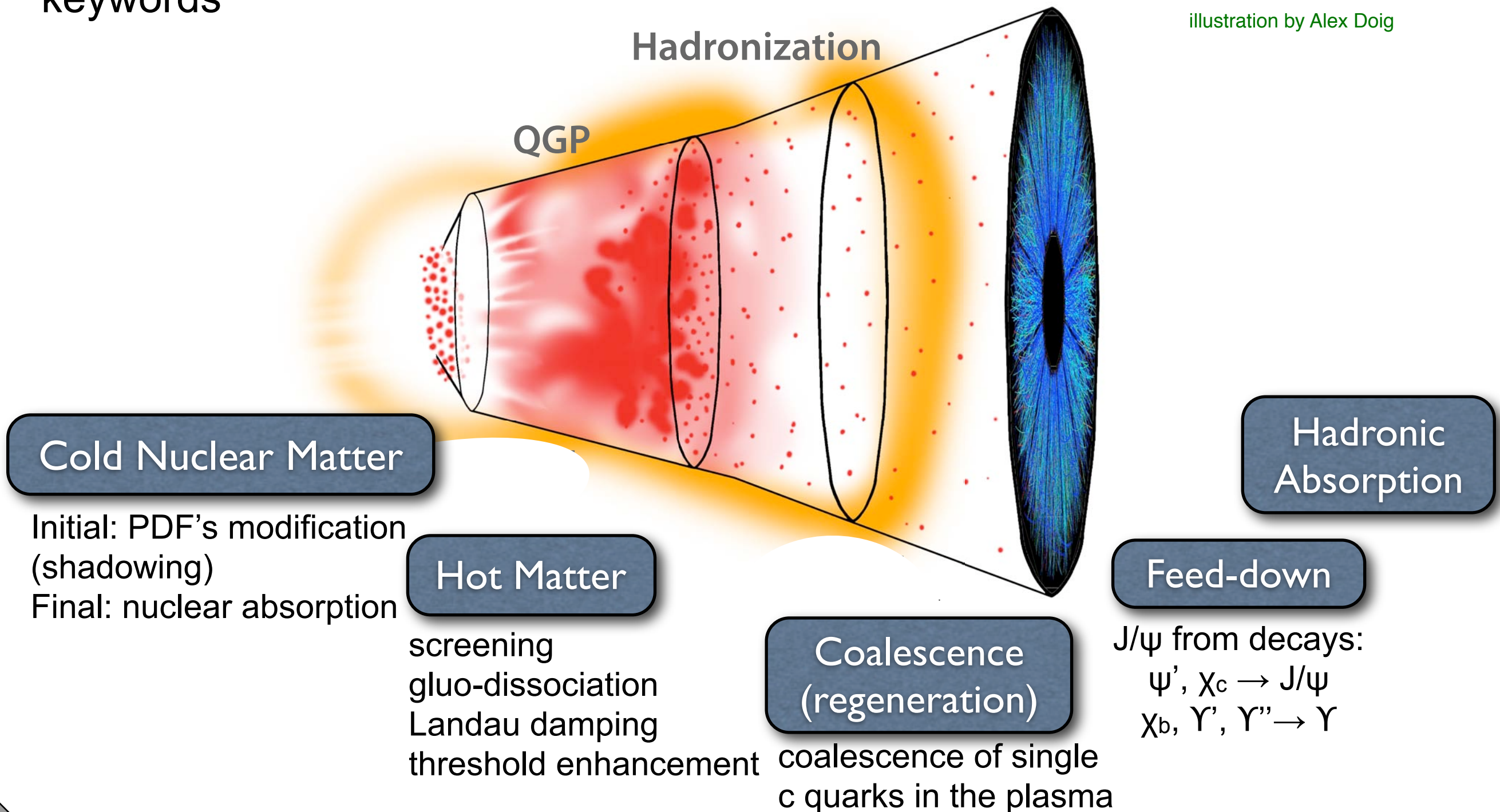
... end up in  
hidden ( $J/\psi$ ) or  
in open (D) charm

# Can we isolate hot effects?



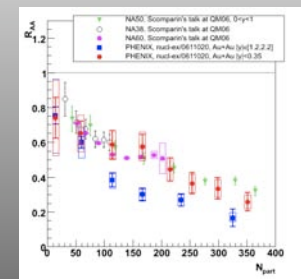
keywords

illustration by Alex Doig



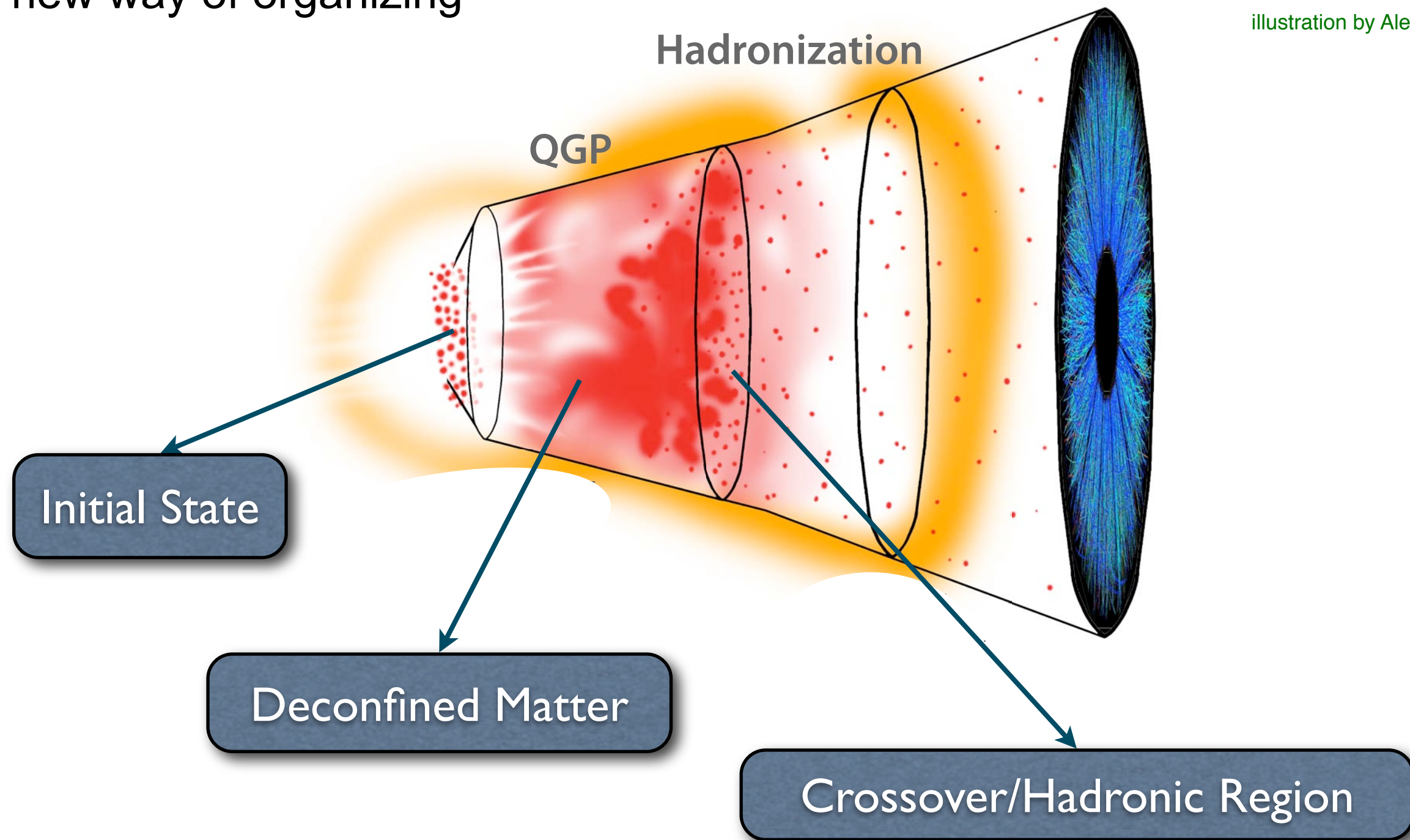


# Can we isolate hot effects?



new way of organizing

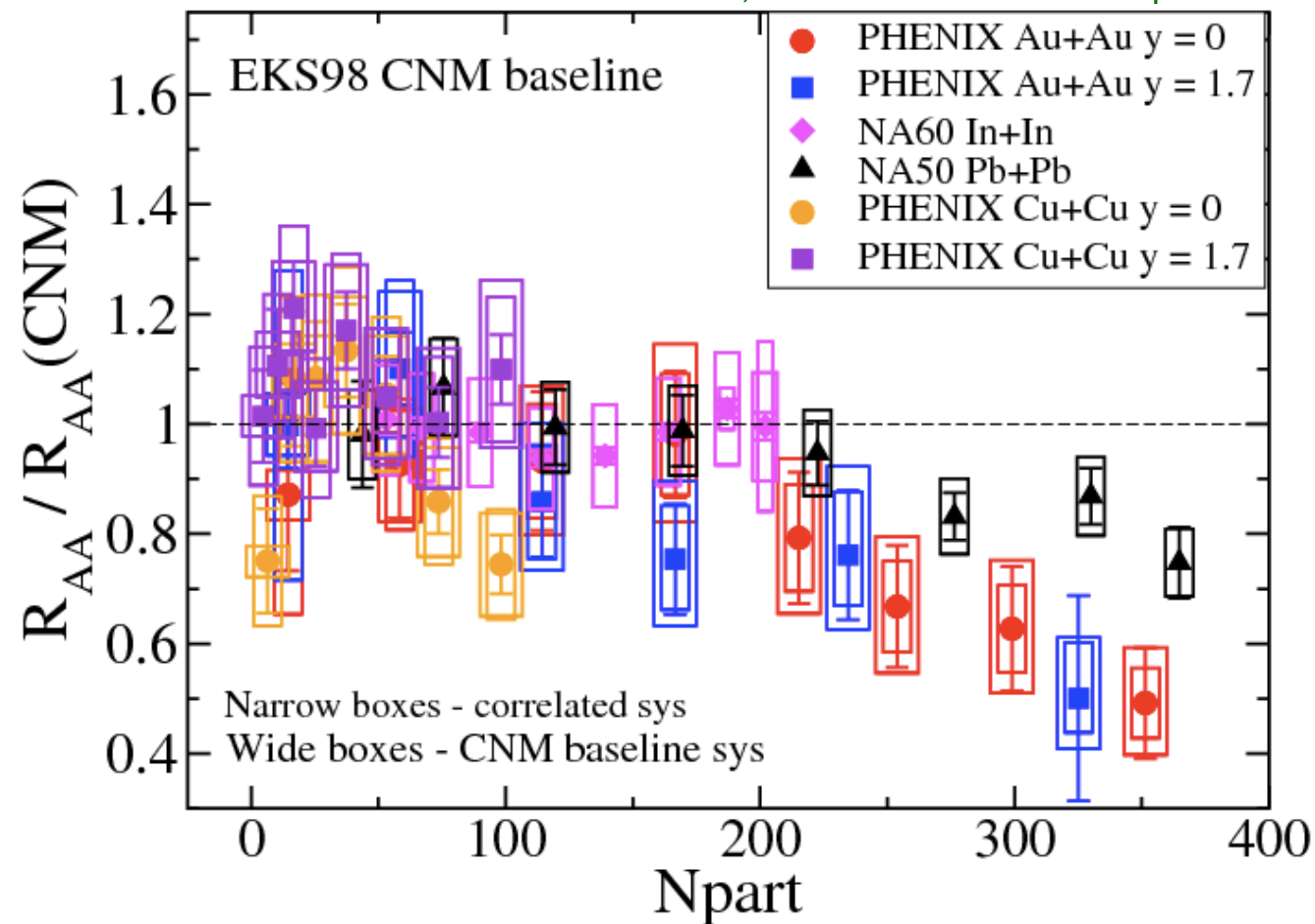
illustration by Alex Doig



# Can we isolate hot effects?

With CNM effects divided out  
“Anomalous suppression”

M. Leitch, INT Quarkonium Workshop 2008



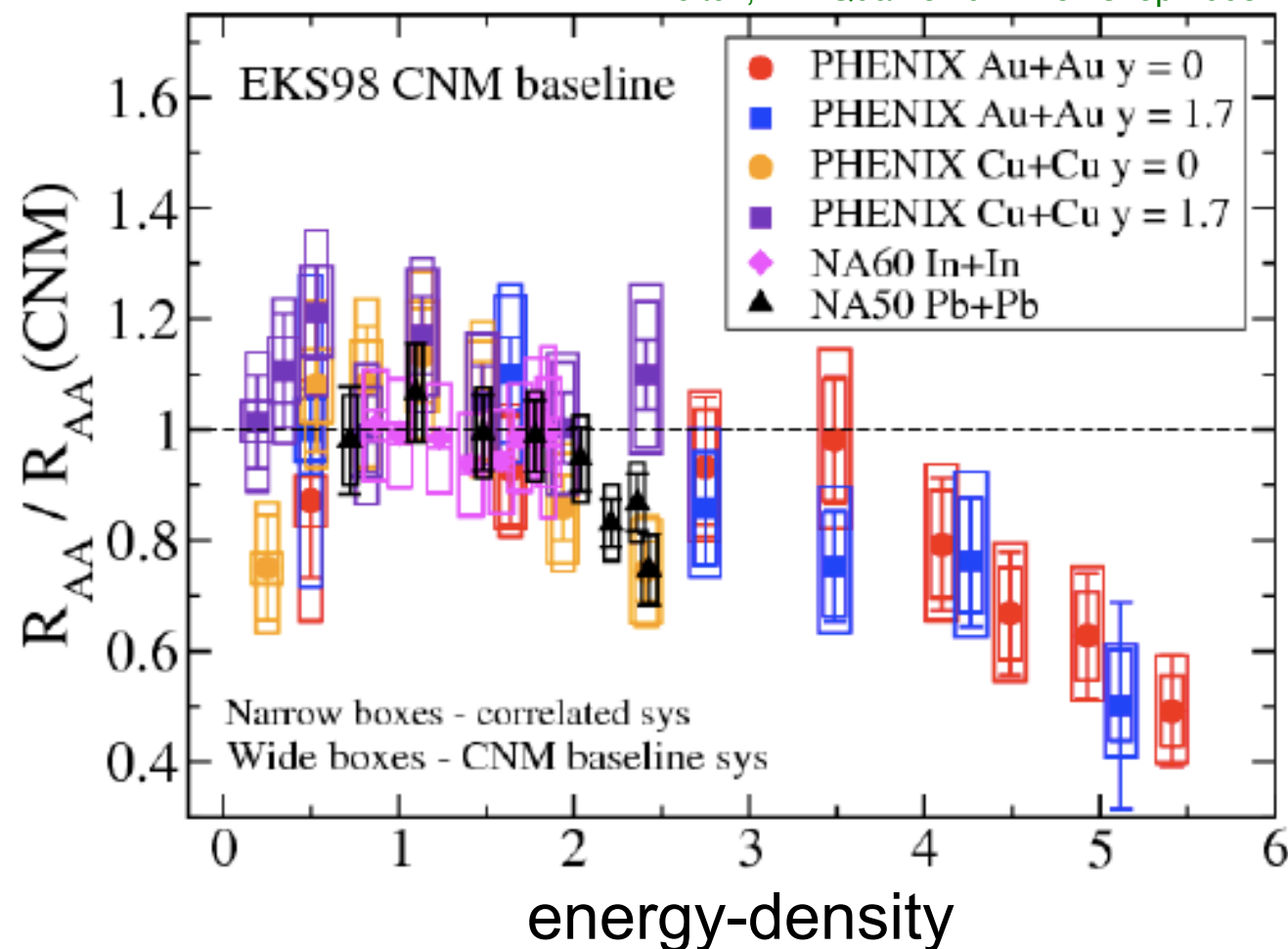
This makes sense if

- all CNM effects are initial state (shadowing), or
- absorption in the crossover region (“mixed phase”) is similar to absorption in nuclear matter

# Can we isolate hot effects?

With CNM effects divided out  
“Anomalous suppression”

M. Leitch, INT Quarkonium Workshop 2008



- SPS described well with hadronic - reaches into crossover region
- RHIC reaches into deconfined region

onset of deconfinement (onset of screening)  $\sim 3.5 \text{ GeV/fm}^3 \sim 200 \text{ MeV}$



# Other Controls

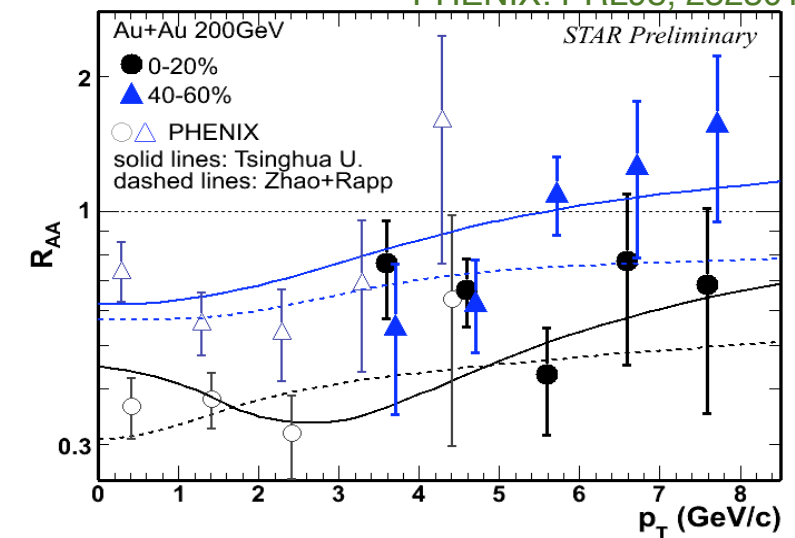
## $R_{AA}$ versus $p_T$

Great for separating the different contributions

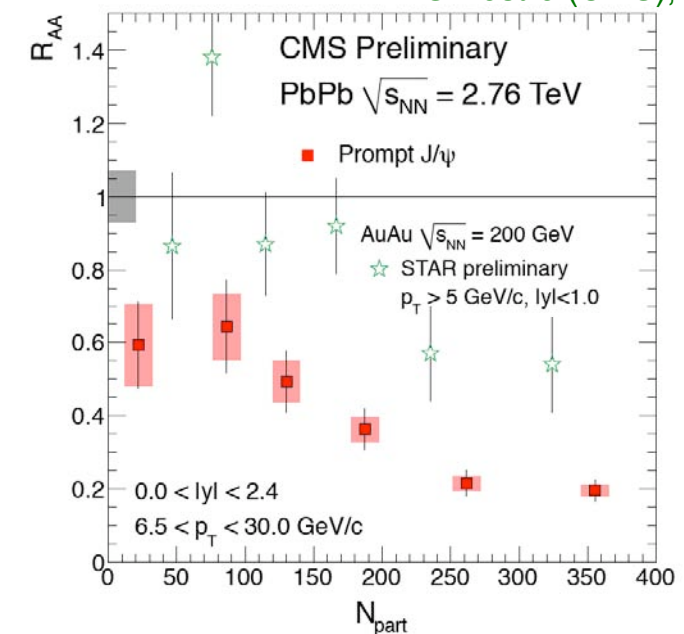
At high  $p_T$

- CNM effects are less important: larger  $x$
- Statistical recombination has little effect
- A suppression at high  $p_T$  would indicate suppression of direct  $J/\psi$  by the hot medium
- If no suppression then  $J/\psi$  forms outside of (or after) the hot plasma. **Formation time!**
- larger high  $p_T$  suppression at LHC: can be from smaller  $x$  and/or longer lifetime

Zebo Tang (STAR), QM 2011  
STAR CuCu: PRC80, 014922(R)  
PHENIX: PRL98, 232301



Silvestre (CMS), QM 2011

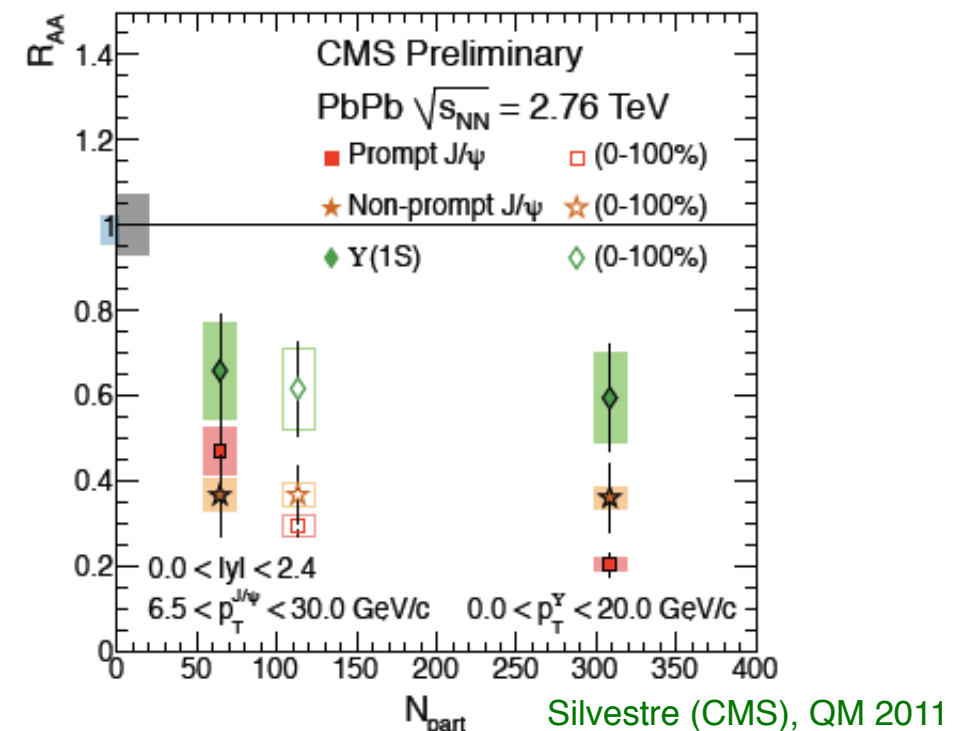
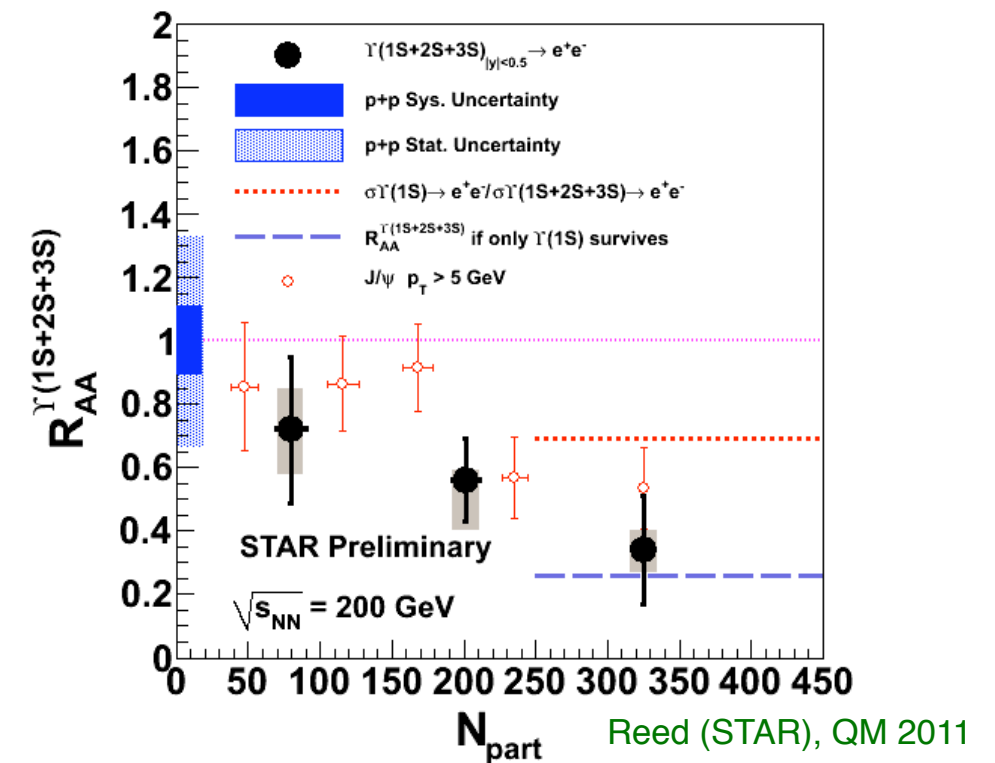


At RHIC: **no suppression** for  $J/\psi$  at high  $p_T$  ( $\sim 5$  GeV) in 200 GeV Cu+Cu and peripheral Au+Au collisions, **but suppression at high  $p_T$**  in central Au  
At LHC suppression persists to higher  $p_T$

# Other Controls

## $\Upsilon$ psilon

- $\Upsilon$  (theoretically) is a much cleaner signal
- Initial state effects not very relevant ( $m_b \gg Q^2$ )
- Absorption is small in the crossover/hadronic region
- No recombination: number of  $b$  and  $b\bar{b}$  is negligible (at RHIC)
- Easy to calculate spectral function, but dynamical modeling harder
- Ground state can survive at RHIC and be suppressed at LHC ?!



# Summary

- **EFT approach:** Framework for systematic studies. The imaginary part of the potential plays a prominent role as a quarkonium dissolution mechanism.
- **Spectral functions determined** - with effects of color screening and dissipation
  - \* Even the most binding potential allowed by lattice QCD it leads to the dissolution of the 1S charmonium & excited bottomonium states by  $T \approx 240$  MeV and of the 1S bottomonium states for  $T \approx 450$  MeV
  - \* Consistent with most recent charmonium spectral functions from lattice
  - \* Different mechanism behind melting of charmonium and bottomonium
  - \* Threshold enhancement has phenomenological consequences
- **Dynamical bridging necessary to make comparison to data** - heavy quark diffusion, strength of  $Q$ - $Q$ bar correlation, lifetime of deconfined medium
- High  $p_T$  - as clear hot matter signal

# The End